

Background to Blackbody Radiation Law

Wien
displacement
law

Derived from
condition that
reversible adiabatic
compression preserves
thermal equilibrium

$$u_\nu(T) = \nu^3 f(\nu/T)$$

↑
Energy density
in frequency
 ν at temperature T

Rayleigh-
Jeans
law
(classical)
1900-1905

$$u_\nu(T) = \frac{8\pi\nu^2}{c^3} kT$$

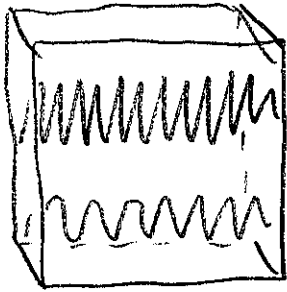
Wien
distribution

$$u_\nu(T) = \frac{8\pi h\nu^3}{c^3} \exp\left(-\frac{h\nu}{kT}\right)$$

Planck
distribution
(fitted to best
data in 1900)

$$u_\nu(T) = \frac{8\pi h\nu^3}{c^3} \frac{1}{\exp\left(\frac{h\nu}{kT}\right) - 1}$$

Direct Pathway



Radiation consists
of oscillators
"normal modes"
of all frequencies
 ν

Density of
oscillators $d\nu$

$$n_\nu = \frac{8\pi\nu^2}{c^3}$$

Classical
analysis.

Each mode has
2 degrees freedom

+

Equipartition

$\frac{1}{2}kT$ energy

per degree

freedom



$$U_\nu(T) = n_\nu kT$$

$$= \frac{8\pi\nu^2}{c^3} kT$$

Quantum
analysis

Each mode is restricted
to energies

$$nE = 0, h\nu, 2h\nu, \dots$$

with probability

proportional to $\exp(-\frac{nE}{kT})$



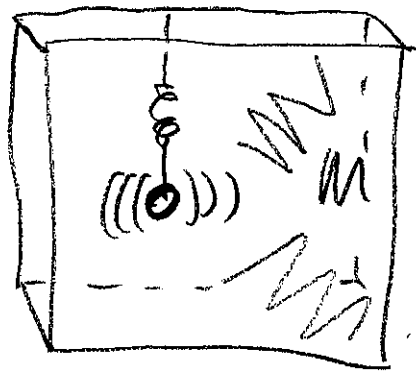
mean
energy
per
mode

$$\bar{E} = \frac{1}{\exp(\frac{h\nu}{kT}) - 1}$$



$$U_\nu(T) = \frac{8\pi\nu^2}{c^3} \cdot \frac{1}{\exp(\frac{h\nu}{kT}) - 1}$$

Indirect Pathway



Radiation is in thermal equilibrium with resonators of 2 degrees of freedom

$$u_\nu(T) = \frac{8\pi\nu^2}{c^3} U$$

Planck. Derived classically before 1900

 ← mean resonator energy at same frequency ν

classical analysis

Equipartition theorem

$$U = 2\left(\frac{1}{2}kT\right) = kT$$

$$u_\nu(T) = \frac{8\pi\nu^2}{c^3} kT$$

Quantum analysis

Resonators limited to energies

$$0, h\nu, 2h\nu, \dots$$

$$U = \frac{1}{\exp\left(\frac{h\nu}{kT}\right) - 1}$$

$$u_\nu(T) = \frac{8\pi\nu^2}{c^3} \frac{1}{\exp\left(\frac{h\nu}{kT}\right) - 1}$$

mean energy \bar{E} of quantized system

Probability
system has
energy

$$P(nE) \propto \exp\left(-\frac{nE}{kT}\right)$$

nE

↑ unit,
usually $h\nu$

mean
energy

$$\bar{E} = \frac{\sum_{n=0}^{\infty} nE \exp\left(-\frac{nE}{kT}\right)}{\sum_{n=0}^{\infty} \exp\left(-\frac{nE}{kT}\right)}$$

$$kT \frac{d}{dE} \left(\underbrace{\sum_{n=0}^{\infty} \exp\left(-\frac{nE}{kT}\right)}_Z \right)$$

$$= kT \cdot \frac{1}{Z} \frac{d}{dE} Z$$

$$Z = \sum_{n=0}^{\infty} \exp\left(-\frac{nE}{kT}\right) = \sum_{n=0}^{\infty} \left(\exp\left(-\frac{E}{kT}\right)\right)^n = \frac{1}{1 - \exp\left(-\frac{E}{kT}\right)}$$

$$kT \frac{d}{dE} Z = kT \frac{1}{\left[1 - \exp\left(-\frac{E}{kT}\right)\right]^2} \cdot \frac{1}{kT} \exp\left(-\frac{E}{kT}\right) = Z^2 \exp\left(-\frac{E}{kT}\right)$$

$$\left[\bar{E} = \frac{Z^2 \exp\left(-\frac{E}{kT}\right)}{Z} = \frac{\exp\left(-\frac{E}{kT}\right)}{1 - \exp\left(-\frac{E}{kT}\right)} = \frac{1}{\exp\left(\frac{E}{kT}\right) - 1} \right]$$