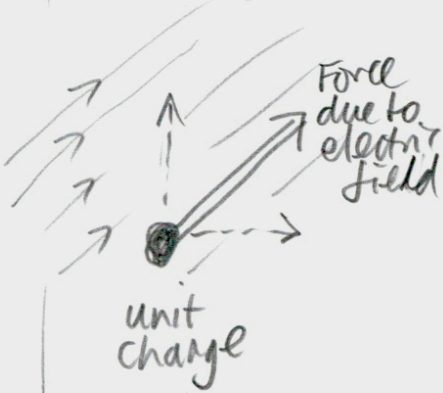


Minimum Maxwell

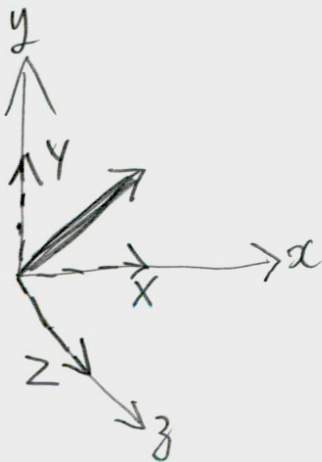
The minimum of Maxwell's equations needed for Einstein's 1905 special relativity paper

Charge free empty space is filled with two fields

Electric field strength

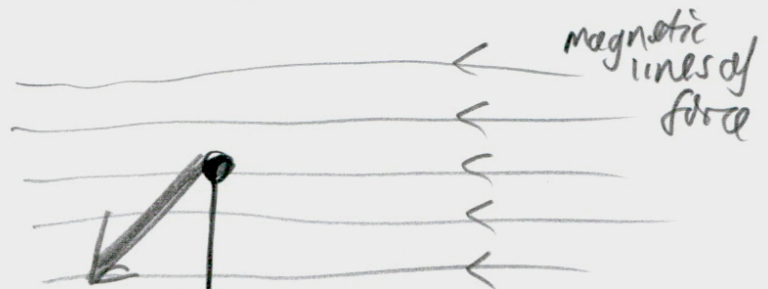


Components of field strength



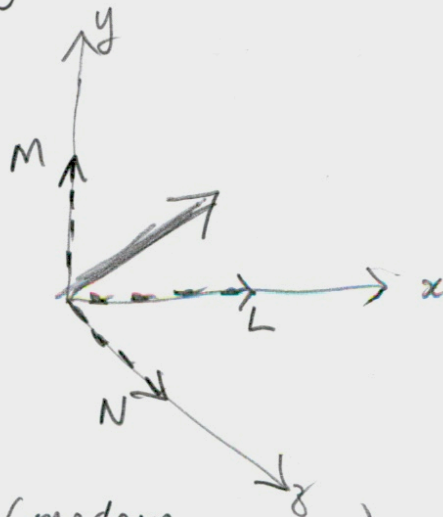
(modern $\vec{E} = (x, y, z)$)

magnetic field strength



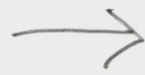
Force = $\vec{v} \times (\text{Field strength})$
 ("Gaussian units")

Components of field strength



(modern $H = (L, M, N)$)

Maxwell's
force
(vector)
equations



Helmholtz fields

$$X(x, y, z), Y(x, y, z), Z(x, y, z)$$

$$L(x, y, z), M(x, y, z), N(x, y, z)$$

are spread over space x, y, z

& how they vary with time t

source free case
No charges

I

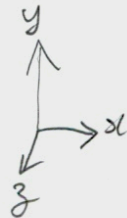
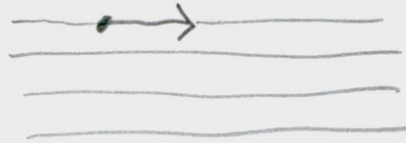
$$\frac{\partial X}{\partial x} + \frac{\partial Y}{\partial y} + \frac{\partial Z}{\partial z} = 0$$

says electric field
lines of force are
conserved

Motivation

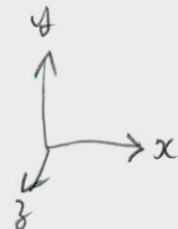
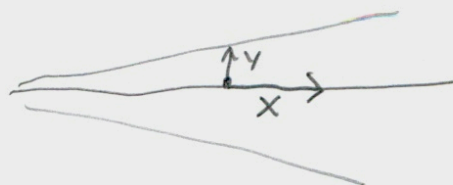
Homogeneous
case

Lines point
in one
direction } x -direction



$$\left. \begin{array}{l} Y=Z=0 \\ \therefore \frac{\partial Y}{\partial y} + \frac{\partial Z}{\partial z} = 0 \end{array} \right\} \Rightarrow \frac{\partial X}{\partial x} = 0 \quad \therefore X = \text{constant}$$

Lines
diverge



$$\text{If } \frac{\partial Y}{\partial y} > 0 \quad \frac{\partial Z}{\partial z} > 0 \Rightarrow \frac{\partial X}{\partial x} < 0 \quad \therefore X \text{ weakens with increasing } x$$

more
precise
grounding

Gauss' theorem

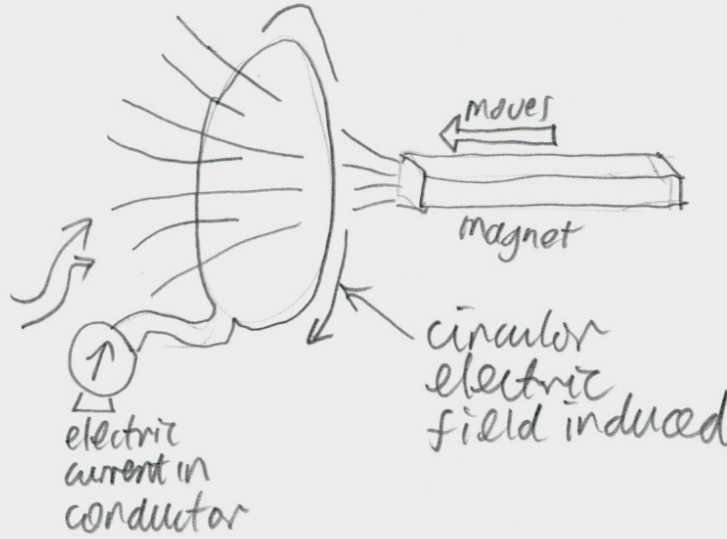
$$\underline{\nabla} \cdot \underline{E} = 0$$

II

$$\frac{\partial L}{\partial x} + \frac{\partial M}{\partial y} + \frac{\partial N}{\partial z} = 0$$

magnetic lines
of force
conserved

Faraday
Induction



Magnetic
field
intensifies
here

"Gaussian
units"

electric
current in
conductor

circular
electric
field induced

III

$$\frac{1}{c} \frac{\partial L}{\partial t} = \frac{\partial Y}{\partial x} - \frac{\partial Z}{\partial y}$$

$$\frac{1}{c} \frac{\partial M}{\partial t} = \frac{\partial Z}{\partial x} - \frac{\partial X}{\partial z}$$

$$\frac{1}{c} \frac{\partial N}{\partial t} = \frac{\partial X}{\partial y} - \frac{\partial Y}{\partial x}$$

More precise
grounding

"Stokes theorem"

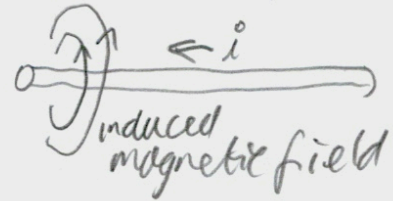
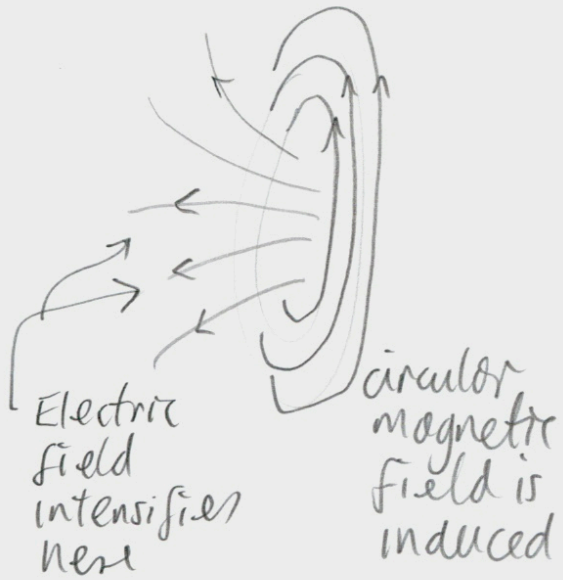
$$-\frac{1}{c} \frac{\partial H}{\partial t} = \nabla \times \underline{E}$$

change
in time
of
magnetic
field

circular
electric
field

Maxwell's
 "Displacement
 currents" } increasing
 electric
 field

induces
 circular
 magnetic
 field
 analogous to
 Oersted/Ampere law



$$\frac{1}{c} \frac{\partial X}{\partial t} = \frac{\partial N}{\partial y} - \frac{\partial M}{\partial z}$$

$$\frac{1}{c} \frac{\partial Y}{\partial t} = \frac{\partial L}{\partial z} - \frac{\partial N}{\partial x}$$

$$\frac{1}{c} \frac{\partial Z}{\partial t} = \frac{\partial M}{\partial x} - \frac{\partial L}{\partial y}$$

change in time of electric field

circular magnetic field

IV

more precise
 grounding
 "Stokes theorem"

$$\frac{1}{c} \frac{\partial \underline{E}}{\partial t} = \nabla \times \underline{H}$$

Einstein's Project

maxwell's equations hold in $K(x, y, z, t)$

Principle of Relativity

maxwell's equations hold in $k(\xi, \eta, \zeta, \tau)$

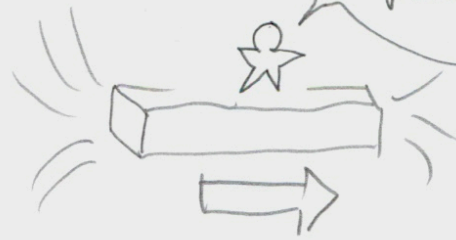
show this is so

x, y, z, t transform by Lorentz transformation

What about field quantities x, y, z ? L, M, N .

magnet and conductor

I see a magnetic field AND an electric field



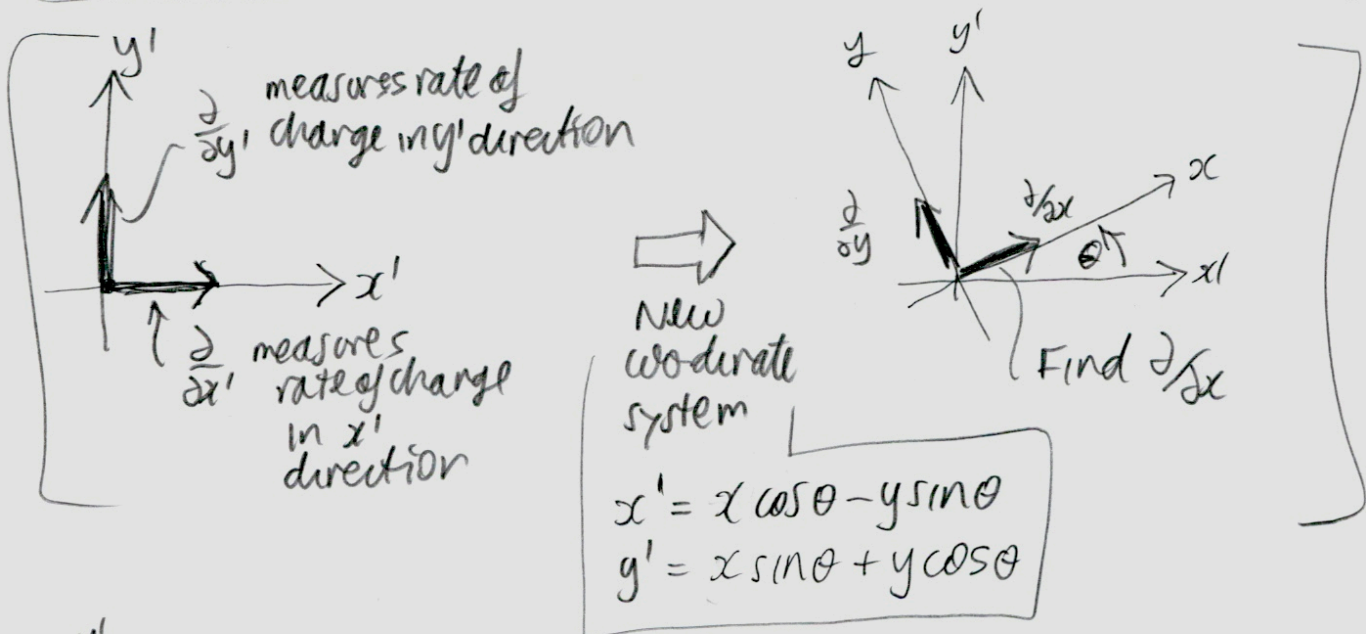
I see a pure magnetic field

Pure magnetic field
 L, M, N
 $x=y=z=0$



magnetic and electric field
 L, M, N
 $x, y, z \neq 0$

Motivating The Chain Rule for Derivative Operators



$$x' = x \cos \theta - y \sin \theta$$

$$y' = x \sin \theta + y \cos \theta$$

$\frac{\partial}{\partial x}$ is fixed by $\frac{\partial}{\partial x'}$, $\frac{\partial}{\partial y'}$

linear dependence

$$\frac{\partial}{\partial x} = (\text{some constant}) \frac{\partial}{\partial x'} + (\text{some constant}) \frac{\partial}{\partial y'}$$

$$\left. \frac{\partial x'}{\partial x} \right|_{y \text{ constant}} = \cos \theta$$

$$\left. \frac{\partial y'}{\partial x} \right|_{y \text{ constant}} = \sin \theta$$

$$\frac{\partial}{\partial x} = \cos \theta \frac{\partial}{\partial x'} + \sin \theta \frac{\partial}{\partial y'}$$

same as rule of vector decomposition

(Derivative operators \sim vectors)

Preparation: Transformations of space, time derivative operators

$$\tau = \beta(t - v/c^2 x) \quad \xi = \beta(x - vt) \quad \eta = y \quad \zeta = z$$

$$\left[\frac{\partial}{\partial t} = \frac{\partial \tau}{\partial t} \cdot \frac{\partial}{\partial \tau} + \frac{\partial \xi}{\partial t} \cdot \frac{\partial}{\partial \xi} + \frac{\partial \eta}{\partial t} \cdot \frac{\partial}{\partial \eta} + \frac{\partial \zeta}{\partial t} \cdot \frac{\partial}{\partial \zeta} = \beta \frac{\partial}{\partial \tau} - \beta v \frac{\partial}{\partial \xi} \right]$$

$$\left[\frac{\partial}{\partial x} = \frac{\partial \tau}{\partial x} \cdot \frac{\partial}{\partial \tau} + \frac{\partial \xi}{\partial x} \cdot \frac{\partial}{\partial \xi} + \frac{\partial \eta}{\partial x} \cdot \frac{\partial}{\partial \eta} + \frac{\partial \zeta}{\partial x} \cdot \frac{\partial}{\partial \zeta} = \beta \frac{\partial}{\partial \xi} - \beta \frac{v}{c^2} \frac{\partial}{\partial \tau} \right]$$

$$\left[\frac{\partial}{\partial y} = \frac{\partial}{\partial \eta} \right]$$

$$\left[\frac{\partial}{\partial z} = \frac{\partial}{\partial \zeta} \right]$$

Inverses $\frac{\partial}{\partial \tau} = \beta \frac{\partial}{\partial t} + \beta \frac{v}{c^2} \frac{\partial}{\partial x} \quad \frac{\partial}{\partial \xi} = \beta \frac{\partial}{\partial x} + \beta v \frac{\partial}{\partial t}$

$$\frac{1}{c} \frac{\partial \underline{E}}{\partial t} = \nabla \times \underline{H} \quad \underline{E} = (X, Y, Z) \quad \underline{H} = (L, M, N)$$

maxwell's equations in
K(x, y, z, t)
transform to K'(ξ, η, ζ, τ)

x-component

$$\frac{1}{c} \frac{\partial X}{\partial t} = \frac{\partial N}{\partial y} - \frac{\partial M}{\partial z}$$



Add $\frac{v}{c} \frac{\partial X}{\partial x}$ to both sides.

where $\frac{v}{c} \frac{\partial X}{\partial x} = -\frac{v}{c} \frac{\partial Y}{\partial y} - \frac{v}{c} \frac{\partial Z}{\partial z}$ since

$$\begin{cases} \nabla \cdot \underline{E} = 0 \\ \frac{\partial X}{\partial x} + \frac{\partial Y}{\partial y} + \frac{\partial Z}{\partial z} = 0 \end{cases}$$

multiply by β

$$\frac{1}{c} \beta \left(\frac{\partial}{\partial t} + v \frac{\partial}{\partial x} \right) X = \beta \frac{\partial (N - \frac{v}{c} Y)}{\partial y} - \beta \frac{\partial (M - \frac{v}{c} Z)}{\partial z}$$



substitute for space, time operators

$$\frac{1}{c} \frac{\partial X}{\partial \tau} = \frac{\partial}{\partial \eta} \beta (N - \frac{v}{c} Y) - \frac{\partial}{\partial \zeta} \beta (M - \frac{v}{c} Z)$$

y-component

$$\frac{1}{c} \frac{\partial Y}{\partial t} = \frac{\partial L}{\partial z} - \frac{\partial N}{\partial x}$$



Add $-\beta^2 \frac{v}{c} \frac{\partial N}{\partial t} + \beta^2 \frac{v}{c} \frac{\partial Y}{\partial x}$ to both sides.

Insert factors $1 = \beta^2 (1 - v^2/c^2)$

$$\frac{1}{c} \beta^2 (1 - \frac{v^2}{c^2}) \frac{\partial Y}{\partial t} - \beta^2 \frac{v}{c^2} \frac{\partial N}{\partial t} + \beta^2 \frac{v}{c} \frac{\partial Y}{\partial x} = \frac{\partial L}{\partial z} - \beta^2 \frac{v}{c^2} \frac{\partial N}{\partial t} + \beta^2 \frac{v}{c} \frac{\partial Y}{\partial x} - \beta^2 (1 - \frac{v^2}{c^2}) \frac{\partial N}{\partial x}$$

$$\frac{1}{c} \beta^2 \frac{\partial Y}{\partial t} - \beta^2 \frac{v}{c^2} \frac{\partial N}{\partial t} + \beta^2 \frac{v}{c} \frac{\partial Y}{\partial x} - \beta^2 \frac{v^2}{c^2} \frac{\partial N}{\partial x} = \frac{\partial L}{\partial z} - \beta^2 \frac{v}{c^2} \frac{\partial N}{\partial t} + \beta^2 \frac{v^2}{c^2} \frac{\partial Y}{\partial t} - \beta^2 \frac{\partial N}{\partial x} + \beta^2 \frac{v}{c} \frac{\partial Y}{\partial x}$$

$$\frac{1}{c} \left(\beta \frac{\partial}{\partial t} + \beta v \frac{\partial}{\partial x} \right) \beta (Y - \frac{v}{c} N) = \frac{\partial L}{\partial z} - \beta \left(\frac{v}{c^2} \frac{\partial}{\partial t} + \frac{\partial}{\partial x} \right) \beta (N - \frac{v}{c} Y)$$



substitute for space, time operators

$$\frac{1}{c} \frac{\partial}{\partial \tau} \beta (Y - \frac{v}{c} N) = \frac{\partial L}{\partial \zeta} - \frac{\partial}{\partial \xi} \beta (N - \frac{v}{c} Y)$$

... etc. for remaining Maxwell equations

see §9 for more details

If Maxwell's equations also hold in k ,

≡ then fields must transform as

$$\begin{aligned}
 \overset{\text{in } k'}{\curvearrowright} X' &= \psi(v) X \overset{\text{in } k}{\curvearrowleft} & L' &= \psi(v) L \\
 Y' &= \psi(v) \beta \left(Y - \frac{v}{c} N \right) & M' &= \psi(v) \beta \left(M + \frac{v}{c} Z \right) \\
 Z' &= \psi(v) \beta \left(Z + \frac{v}{c} M \right) & N' &= \psi(v) \beta \left(N - \frac{v}{c} Y \right)
 \end{aligned}$$

Fix $\psi(v) = \text{since}$

- Transformation forms a group.

$$X \xrightarrow{\text{add } v} X' = \psi(v) X \xrightarrow[\text{subtract } v]{} X'' = \psi(-v) X'$$

But $X'' = X \therefore \psi(-v) \psi(v) X = X$

$\psi(v) \psi(-v) = 1$

- By symmetry

$\psi(v) = \psi(-v)$

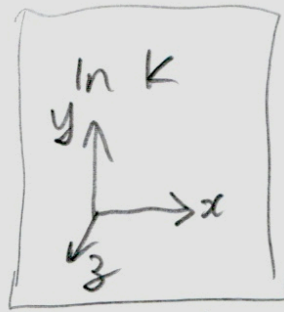
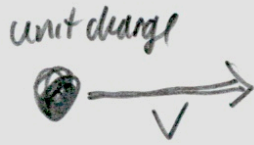
see footnote 4.
 Brief new \underline{E} induced from \underline{H}
 by change frame of reference
 must flip direction if $v \rightarrow -v$.

$$\psi(v) \psi(-v) = (\psi(v))^2 = 1$$

$$\psi(v) = \pm 1$$

↑ choose plus to retain agreement in direction axes
 i.e. roll out $\psi(0) = -1$

"old manner of Expression"



Force = Electric force + magnetic force

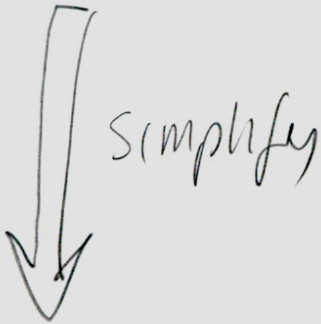
$$\vec{E} = (X, Y, Z)$$

$$\frac{1}{c} \vec{v} \times \vec{H}$$

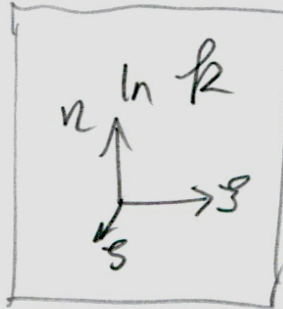
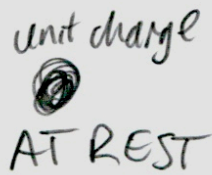
$$= (0, -\frac{v}{c} N, \frac{v}{c} M)$$

use determinant rule

$$\frac{1}{c} \begin{vmatrix} \hat{i} & \hat{j} & \hat{k} \\ v & 0 & 0 \\ L & M & N \end{vmatrix}$$



"New manner of Expression"



Force = Electric force

No magnetic force since charge at rest

Lorentz transform

neglect terms... second and higher in v/c

$$\beta \approx 1$$

$$(X', Y', Z')$$

$$X' = X$$

$$Y' = Y - \frac{v}{c} N$$

$$Z' = Z + \frac{v}{c} M$$

Force in k

Force in K