

Albert Einstein
"On the Inertia
of Energy
Required by the
Relativity
Principle"

Annalen der Physik

23 (1907), pp. 371-384

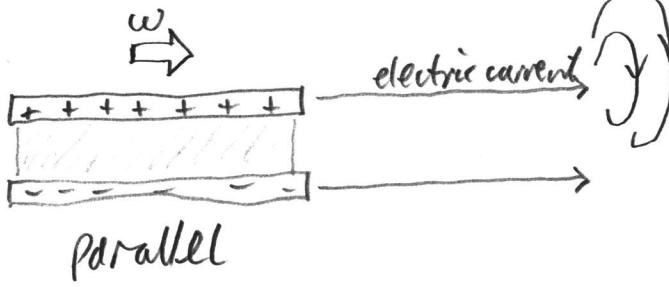
Notes by John D. Norton
September 2015

Introductory puzzle NOT mentioned by Einstein

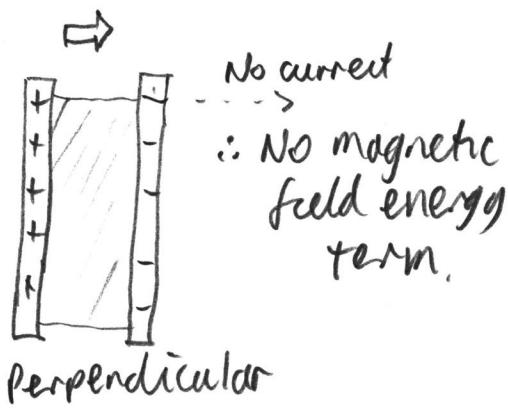
Trotton &
Noble
1903

The other
second order
ether drift
experiment

charged
capacitor
moves at
 ω in
ether



magnetic
field created
with energy
 $\propto (\frac{\omega}{c})^2$ Electro-
static energy
of capacitor



Energy
parallel $>$ Energy
perpendicular

Expect
capacitor to
turn to this
orientation

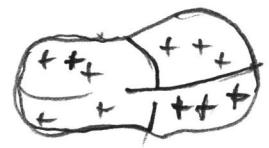
capacitor
suspended
by wire

\Rightarrow no
turning
observed

How do we
now explain
this?

Main Results

1. Kinetic energy of stressed body



$$K = \left[\frac{mc^2}{\sqrt{1-v^2/c^2}} - mc^2 \right] + \text{term due to stress in the body}$$

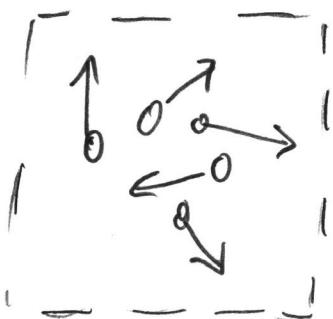
↑
familiar term

$$\Delta E = \frac{v^2}{c^2} \frac{1}{\sqrt{1-v^2/c^2}} (\text{Rest volume}) \begin{pmatrix} \text{Normal} \\ \text{stress in} \\ \text{direction} \\ \text{motion} \end{pmatrix}$$

Energy NOT due to elastic deformation

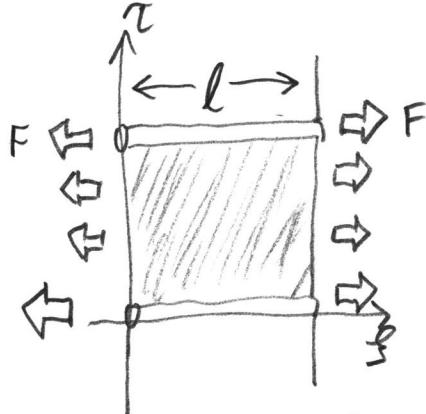
2. Rigid bodies require signals to propagate at speeds $>$ speed light. [(novel?) energy argument]
 Possibility of signals arriving before they leave
 "impossible"

3. Collection of moving masses : Kinetic energy of each mass contributes to rest mass of the total system

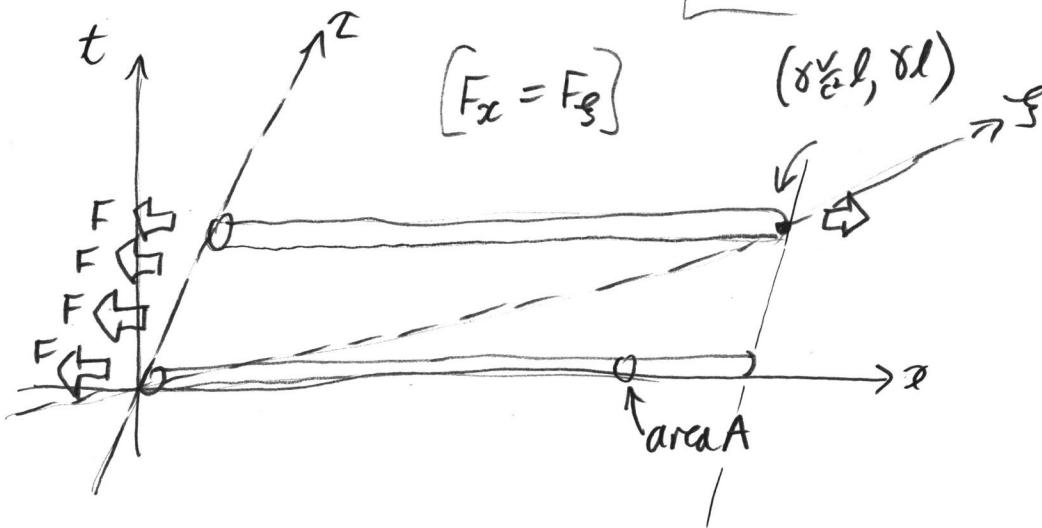
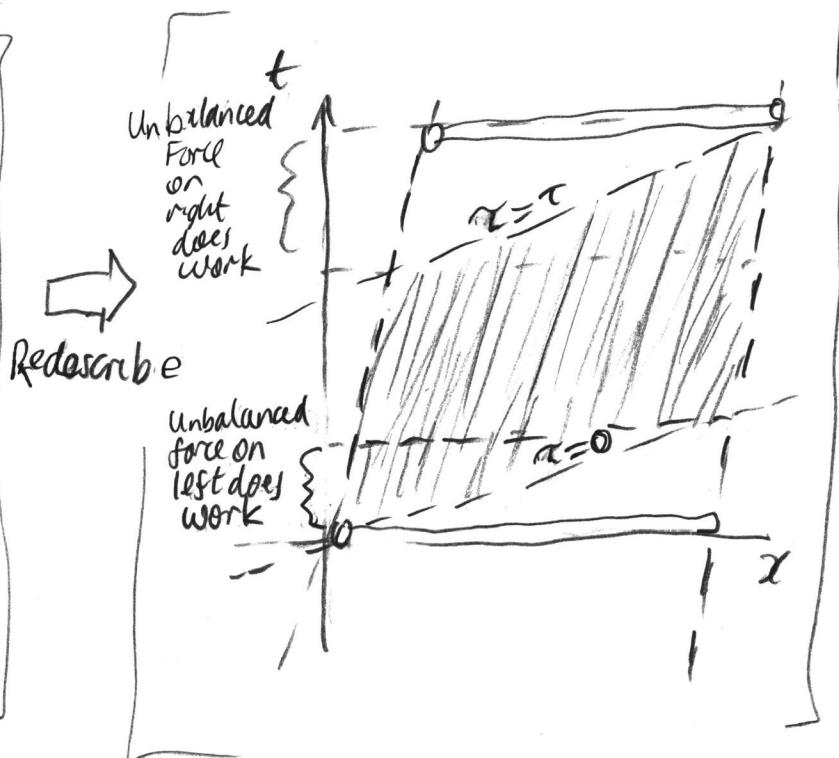


Simplified version of result 1

In $\delta_k(\xi, \eta, \zeta, \tau)$, rod at rest



Rod in tension between times $\tau=0, \tau=\tau$ ONLY



$$\begin{aligned} \tau &= 0 \\ l &= l \\ \downarrow \\ t &= \gamma(\tau + \sqrt{c^2}) \\ &= \gamma \sqrt{c^2} l \\ x &= \gamma(\xi + \sqrt{c^2}) \\ &= \gamma l \\ \gamma &= \sqrt{1 - v/c} \end{aligned}$$

$$\left(\begin{array}{l} \text{Work done} \\ \text{on rod} \\ \text{during} \\ \text{time} \\ \gamma \sqrt{c^2} l \\ \text{of unbalanced} \\ \text{force} \end{array} \right) = -F \cdot V \cdot \gamma \frac{v}{c^2} l$$

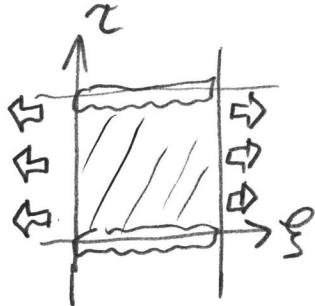
$$= -\frac{F}{A} \cdot A l \cdot \frac{V^2}{c^2} \cdot \gamma$$

$$= \left(\begin{array}{l} \text{(Normal)} \\ \text{stress} \end{array} \right) \left(\begin{array}{l} \text{(Rest} \\ \text{volume)} \end{array} \right) \frac{V^2/c^2}{\sqrt{1 - V^2/c^2}}$$

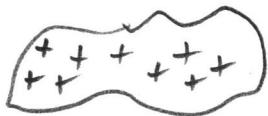
$$\begin{aligned} &= \Delta E \\ &\text{energy} \\ &\text{change in} \\ &\text{rod} \end{aligned}$$

Einstein's more general case

$K(\xi, \eta, \zeta, t)$ Rest frame of electrically charged body



External field present only
in $t=0$ to $t=T$



charge density p'

Net force in
 x -direction

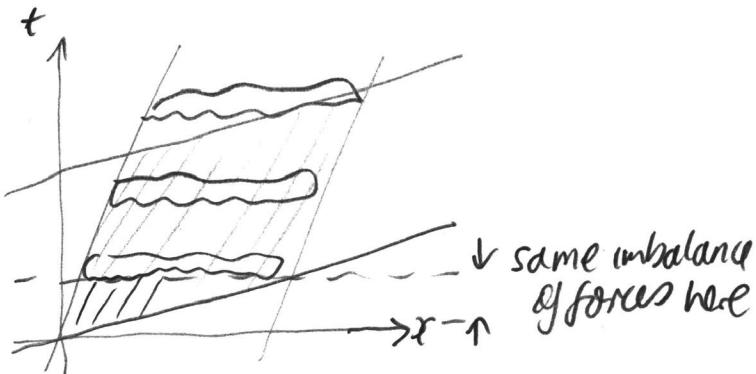
$$= \int x' \frac{p'}{4\pi} d\xi d\eta d\zeta$$

$$= 0$$

equilibrium

x'
 x -component
of E

Transform to $K(x, y, z, t)$
where body moves at
 v in x -direction



compute energy during brief period of unbalanced force:

$$\Delta E = - \frac{(v/c)^2}{\sqrt{1-v^2/c^2}} \sum \xi K_E$$

\uparrow proper position in body \uparrow K_E \uparrow ξ component of force on rest volume element $d\xi d\eta d\zeta$

$$= \frac{x'_i p'_i}{4\pi} d\xi d\eta d\zeta$$

sum over all volume elements

(Long,
messy
integration)

Long Term Significance

Energy density

represented as O-O component

O-O

component

(STRESS-energy tensor)

Lorentz transformation
mixes energy density & stresses

$T =$

$$\begin{bmatrix} \text{Energy density} & | & \text{Energy flux} = \text{momentum density} \\ \hline & | & | \\ \text{Energy flux} & | & \text{Normal stress} & \text{Shear stress} \\ " & | & | & | \\ \text{momentum density} & | & \text{Normal stress} & \text{Normal stress} \\ | & | & | & | \\ \text{shear stress} & & & \text{Normal stress} \end{bmatrix}$$

Rest frame
fluid under
pressure

boost
in
 x -direction \Rightarrow

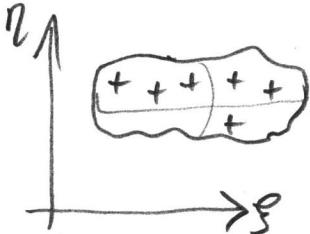
** Term
Linear in
rest energy density,
pressure

$$T = \begin{bmatrix} \text{rest energy density} & | & 0 \\ | & | & | \\ 0 & | & \text{Press} \\ | & | & | \\ 0 & | & \text{Press} \\ | & | & | \\ & | & \text{Press} \end{bmatrix}$$

$$T = \begin{bmatrix} \text{Energy density} & | & | & | \\ | & | & | & | \\ | & | & | & | \\ | & | & | & | \end{bmatrix}$$

Apparent Anisotropy of e-m self energy of charged body

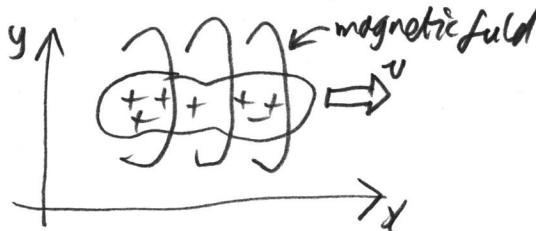
Body at rest in $k(\xi, \eta, \zeta, t)$



Electric field $(x', y', z') \neq 0$

magnetic field $(L', M', N') = 0$
Electrostatic system

Body moves at v in $+x$ direction in $K(x, y, z, t)$



From 1905 special relativity paper

$$\begin{aligned} x &= x' & l &= l' \\ y &= \beta(y' + \frac{v}{c}N') & m &= \beta(m' - \frac{v}{c}z') \\ z &= \beta(z' - \frac{v}{c}m') & n &= \beta(n' + \frac{v}{c}y') \end{aligned}$$

$$\beta = \frac{1}{\sqrt{1 - \frac{v^2}{c^2}}}$$

Total

$$\text{electromagnetic self-energy in } K = E_e = \frac{1}{8\pi} \int x^2 + y^2 + z^2 + l^2 + m^2 + n^2 dx dy dz$$

Substitute k frame values in formula

$$dx = \sqrt{\eta^2 ds} = \frac{1}{\beta} d\xi \quad x = x' \quad l = l'$$

$$y = \beta(y' + \frac{v}{c}N') = \beta y' \quad \left. \right\} y^2 + N^2 = \beta^2 (y'^2 + \frac{v^2}{c^2} y'^2)$$

$$N = \beta(N' + \frac{v}{c}y') = \beta \frac{v}{c} y' \quad \left. \right\} = \frac{1 + (\frac{v}{c})^2}{1 - (\frac{v}{c})^2} y'^2$$

similar

$$z^2 + m^2 = \frac{1 + (\frac{v}{c})^2}{1 - (\frac{v}{c})^2} z'^2$$

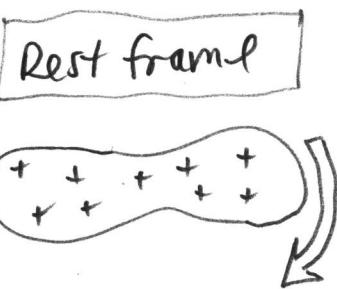
$$E_e = \frac{1}{8\pi} \left\{ \frac{1}{\beta} \left[x'^2 + \frac{1 + (\frac{v}{c})^2}{1 - (\frac{v}{c})^2} (y'^2 + z'^2) \right] d\xi d\eta d\zeta \right\}$$

Self-energy of moving charged body varies with orientation w.r.t motion

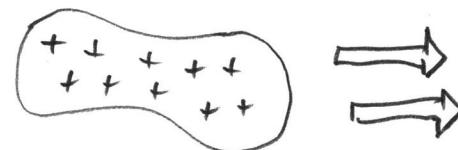
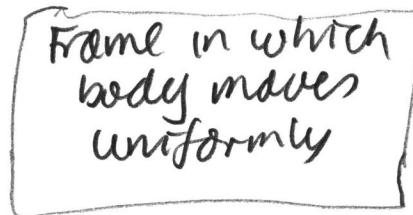
JDW (Einsteins not say this)
Effect slight in Trouton-Noble experiment

Einstein's Worry

contradiction with
principle of relativity



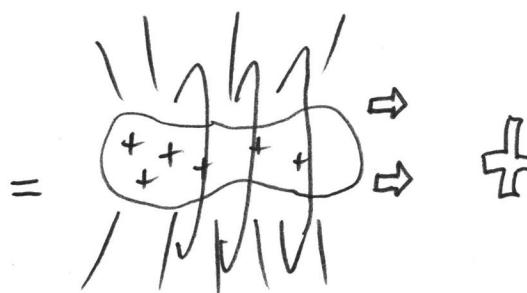
charged
body can
rotate "infinitely
slowly"



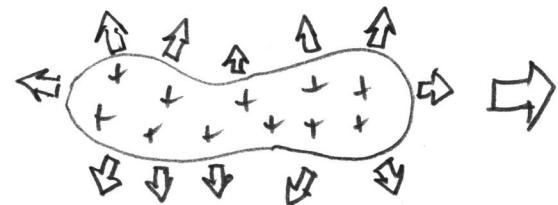
charged body cannot
rotate infinitely slowly
but will settle in
lowest energy state

Einstein's Solution |

Energy of moving body due to electric charge



Energy of its electromagnetic field



Energy associated with stresses induced by charges

$$E_e =$$

$$\frac{1}{8\pi} \int \frac{1}{\beta} \left[x'^2 + \frac{1 + (\frac{v}{c})^2}{1 - (\frac{v}{c})^2} (y'^2 + z'^2) \right] d\sigma dy dz$$

Anisotropic

$$\Delta E = -\frac{v}{c^2} \beta \frac{1}{4\pi} \int \rho' d\sigma dy dz$$

$\uparrow \frac{\partial x'}{\partial \sigma} + \frac{\partial y'}{\partial \eta} + \frac{\partial z'}{\partial \zeta}$

Anisotropic

sum is isotropic!

$$E_e + \Delta E = \underbrace{\left[\frac{1}{8\pi} \int (x'^2 + y'^2 + z'^2) d\sigma dy dz \right]}_{\text{Rest energy of electrostatic field } E_s} \frac{1}{\sqrt{1 - v^2/c^2}}$$

Rest energy of electrostatic field E_s

$$= \frac{\left(\frac{E_s}{c^2} \right) c^2}{\sqrt{1 - v^2/c^2}}$$

Electric charge add mass
 E_s/c^2
that dilates with speed isotropically like ordinary mass

Evaluating $E_e + \Delta E$

$$\Delta E = -\frac{v^2}{c^2} \beta \frac{1}{4\pi} \int_{\text{all space}} \xi x' \left(\frac{\partial x'}{\partial \xi} + \frac{\partial y'}{\partial \eta} + \frac{\partial z'}{\partial \zeta} \right) d\xi d\eta d\zeta$$

Integrate by parts



$$\int_{\text{all space}} \left(\frac{\partial}{\partial \xi} \xi x'^2 + \frac{\partial}{\partial \eta} \xi x' y' + \frac{\partial}{\partial \zeta} \xi x' z' \right) - x'^2 - x' \xi \frac{\partial x'}{\partial \xi} - y' \xi \frac{\partial x'}{\partial \eta} - z' \xi \frac{\partial x'}{\partial \zeta} d\xi d\eta d\zeta$$

By Gauss theorem equal to
an integral over an
infinitely distant surface

of vector

$$(x', y', z')$$

$\sim 1/\text{radius}^4$ or $1/\text{radius}^3$

$$-y' \xi \frac{\partial^2 \phi}{\partial \eta \partial \xi}$$

$$= -y' \xi \frac{\partial^2 \phi}{\partial \xi \partial \eta}$$

$$= y' \xi \frac{\partial y'}{\partial \xi} \quad z' \xi \frac{\partial z'}{\partial \xi}$$

\therefore vanishes

$$(x', y', z') = \left(-\frac{\partial \phi}{\partial \xi}, -\frac{\partial \phi}{\partial \eta}, -\frac{\partial \phi}{\partial \zeta} \right)$$

$$= \int_{\text{all space}} -x'^2 - \xi x' \frac{\partial x'}{\partial \xi} - \xi y' \frac{\partial y'}{\partial \xi} - \xi z' \frac{\partial z'}{\partial \xi} d\xi d\eta d\zeta$$

$$-\frac{1}{2} \xi \frac{\partial}{\partial \xi} (x'^2 + y'^2 + z'^2)$$

Integrate by parts again

OVER

$$= \int_{\text{all space}} -x'^2 - \frac{1}{2} \frac{\partial^2}{\partial s^2} \left(\xi(x'^2 + y'^2 + z'^2) + \frac{1}{2} (x'^2 + y'^2 + z'^2) \frac{\partial \xi}{\partial s} \right) d\xi d\eta ds$$

$$\int_{\text{all space}} \frac{\partial^2}{\partial s^2} \left(\xi(x'(s, \eta, s) + \dots) \right) d\xi d\eta ds$$

$$= \int_{\substack{\text{all} \\ \eta, s}} \xi(x'^2(+\infty, \eta, s) + \dots) - \xi(x'^2(-\infty, \eta, s)) d\eta ds$$

$$= \int_{\substack{\text{all} \\ \eta, s}} 0 d\eta ds = 0$$

$$= -\frac{1}{2} \int_{\text{all space}} x'^2 - y'^2 - z'^2 d\xi d\eta ds$$

Combining:

$$\Delta E = \frac{v^2}{c^2} \beta \frac{1}{8\pi} \int_{\text{all space}} x'^2 - y'^2 - z'^2 d\xi d\eta ds$$

$$E_e + \Delta E = \frac{1}{8\pi} \cdot \frac{1}{\beta} \left(x'^2 + \left[\frac{1+(v_e)^2}{1-(v_e)^2} \right] (y'^2 + z'^2) \right) d\zeta d\eta d\zeta$$

$$+ \frac{v_e^2}{c^2} \frac{1}{8\pi} \beta \left(x'^2 - y'^2 - z'^2 \right) d\zeta d\eta d\zeta$$

$$\frac{1}{8\pi} \beta \left(\frac{1}{\beta^2} \left(x'^2 + \left[\frac{1+(v_e)^2}{1-(v_e)^2} \right] (y'^2 + z'^2) \right) d\zeta d\eta d\zeta \right) \Delta E$$

$$(1 - \frac{v_e^2}{c^2})$$

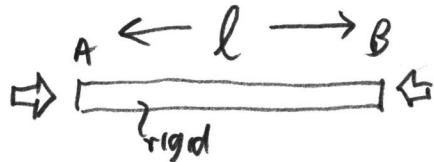
$$(1 - \frac{v_e^2}{c^2}) X'^2 + (1 + \frac{v_e^2}{c^2}) [Y'^2 + Z'^2] = (X'^2 + Y'^2 + Z'^2) - \frac{v_e^2}{c^2} (X'^2 - Y'^2 - Z'^2)$$

$$\therefore E_e + \Delta E = \beta \frac{1}{8\pi} \int x'^2 + y'^2 + z'^2 d\zeta d\eta d\zeta$$

$$E_s$$

Impenetrability of rigid rods/signals at $v > c$

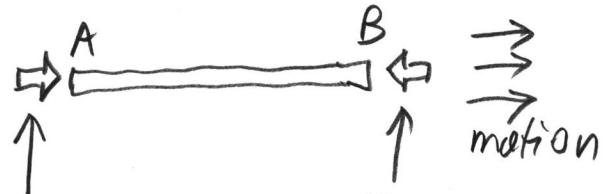
In rest frame of rod



Forces act for very short time at ends of rod

\Downarrow
No change in rest state rod

In frame in which rod moves at v



Force acts EARLIER by time $\frac{v}{c^2} \beta l$

Force acts later

Rod's motion unchanged

Force at A does work

\Rightarrow supplies energy to rod

BUT

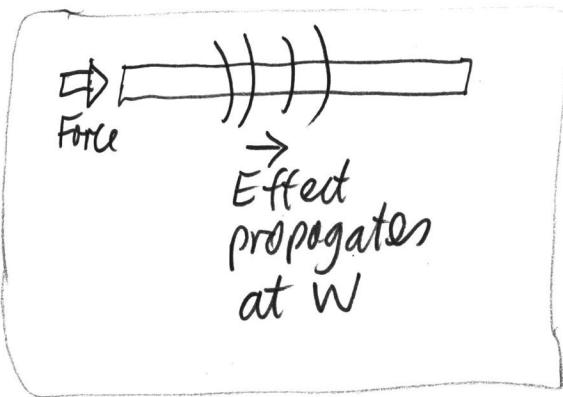
Rod's state does not change in any way.

\therefore Energy of rod \neq function (state of rod)

Hence rigidity untenable

(Novel?) Dynamical argument

Kinematic causal argument



view in frame that moves at $+v$

$$\frac{w'}{v} = \frac{w-v}{1-\frac{wv}{c^2}}$$

signal moves
at w'
covers
distance l
in time T

$$w' = \frac{l}{T} = \frac{w-v}{1-\frac{wv}{c^2}}$$

$$\therefore T = \frac{l(1-\frac{wv}{c^2})}{w-v}$$

↓ if $w > c$

"effect... precedes
cause (accompanied
by act of will, for
example)"

"... does not
contain a contradiction
from a purely logical
point of view..."

conflicts so absolutely
with the character
of all experience, that
the impossibility of the
assumption ... proved..."

$$= \frac{l}{w-v} \cdot \left(1 - \frac{w}{c} \cdot \frac{v}{c}\right)$$

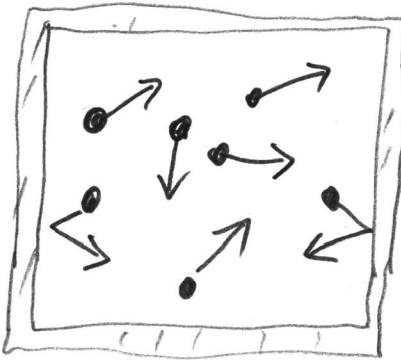
↑
greater
than
1

$v < c$ can be
brought close enough
to c so that
 $\left(1 - \frac{w}{c} \cdot \frac{v}{c}\right) < 0$

$$\boxed{T < 0}$$

Energy of a system consisting of mass points moving force-free

The obvious question.



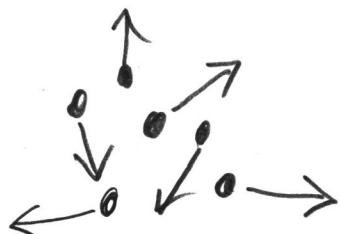
Each molecule of an ideal gas with rest mass m has total energy $mc^2/\sqrt{1-v^2/c^2}$

Does this ^{total} energy contribute to the rest mass of the gas system?

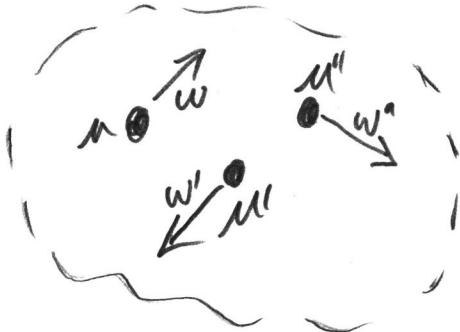
Not quite the question Einstein answers.

Gas pressure \Rightarrow stress in chamber wall \Rightarrow Energy term ΔE
if gas + chamber moves as a whole.

Einstein just considers system of moving masses WITHOUT containment



$K(\xi, \eta, \zeta)$ is the rest frame in the sense that total momentum = 0



$$\sum \frac{mw_1}{\sqrt{1-(w_1/c)^2}} = 0$$

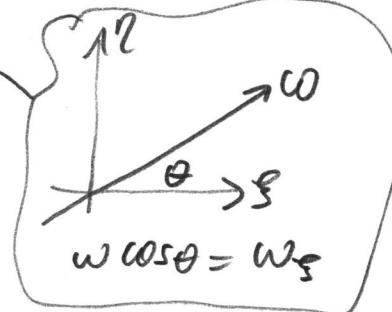
$$\sum \frac{mw_2}{\sqrt{1-(w_2/c)^2}} = 0$$

$$\sum \frac{mw_3}{\sqrt{1-(w_3/c)^2}} = 0$$

In $K(x, y, z)$... K moves at v in +x direction:

Energy E of each mass $E = mc^2 \frac{1+w_1 \cos \theta}{c^2}$
 ("easily determined")

$$\sqrt{1-v^2/c^2} \sqrt{1-w^2/c^2}$$



$$\therefore \text{Total energy } E = \sum_{\text{all masses}} E = \frac{1}{\sqrt{1-v^2/c^2}} \left[\sum mc^2 \cdot \frac{1}{\sqrt{1-w^2/c^2}} \right]$$

(Total energy/c)
 (in (momentum=0) frame) = (rest mass of total system)

$$+ \frac{v}{\sqrt{1-v^2/c^2}} \left[\sum \frac{mw \cos \theta}{\sqrt{1-w^2/c^2}} \right]$$

0 by momentum condition

$$E = \left[\frac{\sum \frac{mc^2}{\sqrt{1-w^2/c^2}}}{c^2} \right] c^2 \quad ***$$