

De Broglie's "Theorem of Phase Agreement"

Periodic phenomena
fixed to moving
particle ...
frequency

$$\nu_1 = \frac{m_0 c^2}{h} \sqrt{1 - \beta^2}$$

$$= \nu_0 \sqrt{1 - \beta^2}$$

v.r.t.
observer at rest

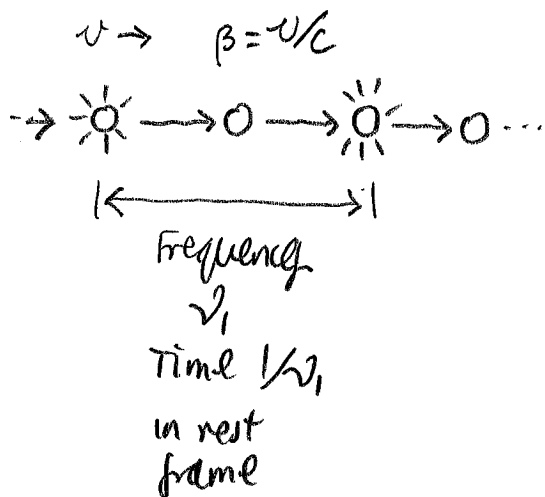
appears
to
this
observer
always
in phase
with

a wave of frequency

$$\nu = \frac{m_0 c^2}{h} \frac{1}{\sqrt{1 - \beta^2}}$$

$$= \nu_0 / \sqrt{1 - \beta^2}$$

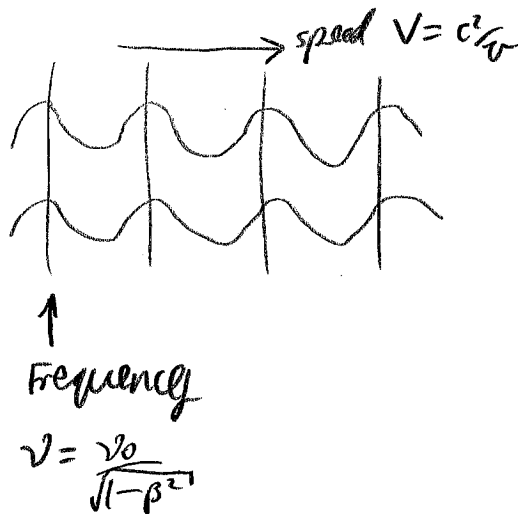
propagating in the
direction of motion
of the moving particle
at ... $V = c/\beta = c^2/\nu$



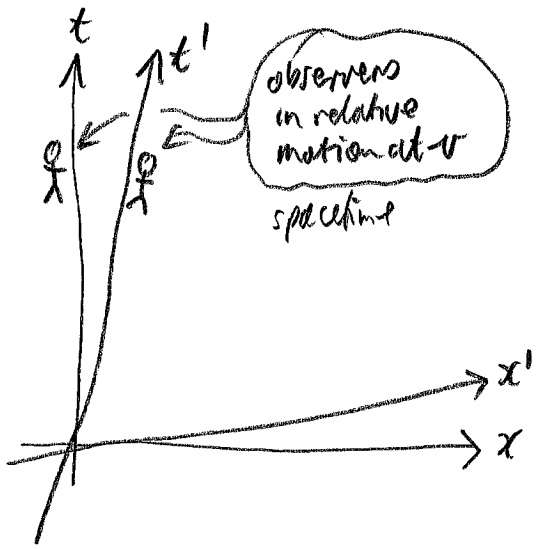
$$\nu_1 = \sqrt{1 - \beta^2} \nu_0$$

$\nu_0 =$ proper
frequency
of object

$$h\nu_0 = m_0 c^2$$



Lorentz transformation



$$t' = \gamma (t - v/c^2 x)$$

$$x' = \gamma (x - vt)$$

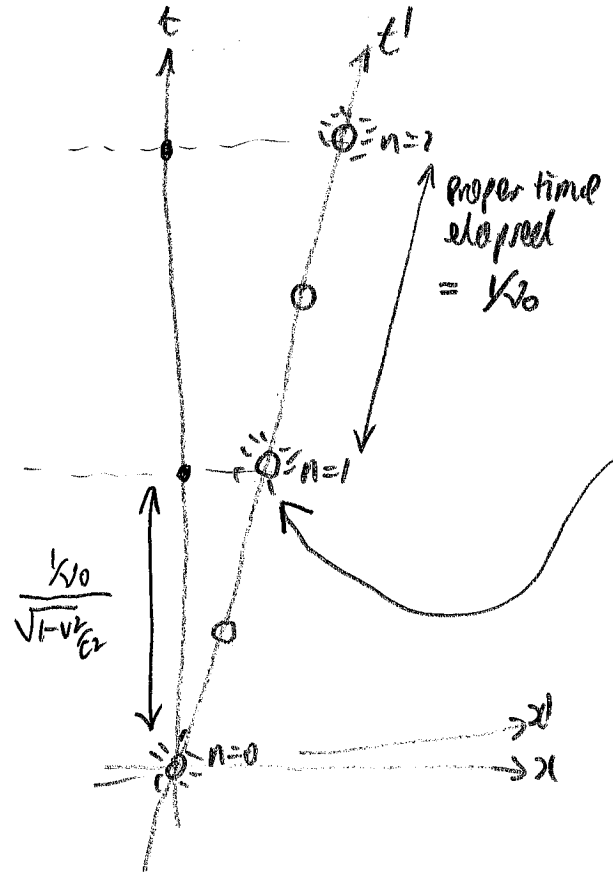
$$\gamma = \frac{1}{\sqrt{1 - v^2/c^2}}$$

$$t = \gamma (t' + v/c^2 x')$$

$$x = \gamma (x' + vt')$$

Particle with proper frequency $\nu_0 = \frac{m_0 c^2}{h}$ at rest in primed frame

particle moves at
 $x = vt$
 $t = \frac{1}{v_0} x$



$n=1$ flash happens at
 $(x', t') = (0, 1/\nu_0)$

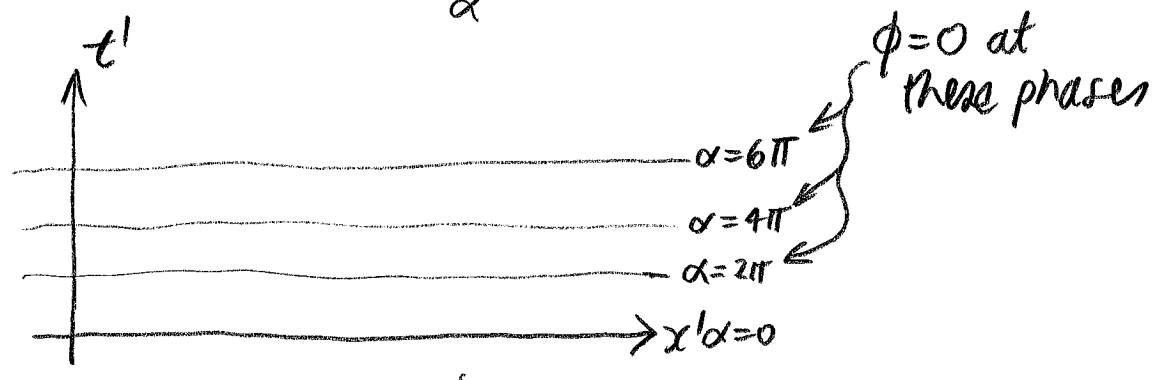
↓ transform to unprimed frame

$$t = \gamma (t' + v/c^2 x') = \gamma t' = \frac{1/\nu_0}{\sqrt{1 - v^2/c^2}}$$

Hence frequency in unprimed frame is
 $\nu_1 = \nu_0 \sqrt{1 - v^2/c^2}$

Plane wave with proper frequency ν_0 for a co-moving observer

$$\phi(x', t') = \sin(\underbrace{2\pi\nu_0 t'}_{\alpha})$$



Transform to unprimed frame

$$\begin{aligned} \phi(x, t) &= \sin(2\pi\nu_0 \delta(t - \frac{v}{c}x)) \\ &= \sin(2\pi(\underbrace{\delta\nu_0}_{\text{Frequency}} t - \underbrace{\frac{\delta\nu_0 v}{c}}_{\text{wavelength}} x)) \end{aligned}$$

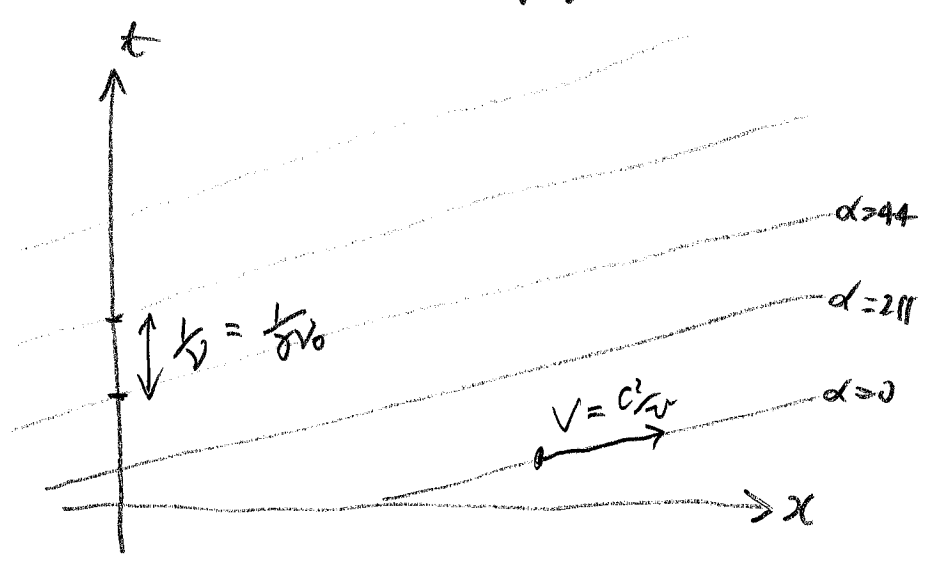
$\nu = \delta\nu_0 = \frac{\nu_0}{\sqrt{1 - v^2/c^2}}$

Phase velocity = velocity of a point of constant phase

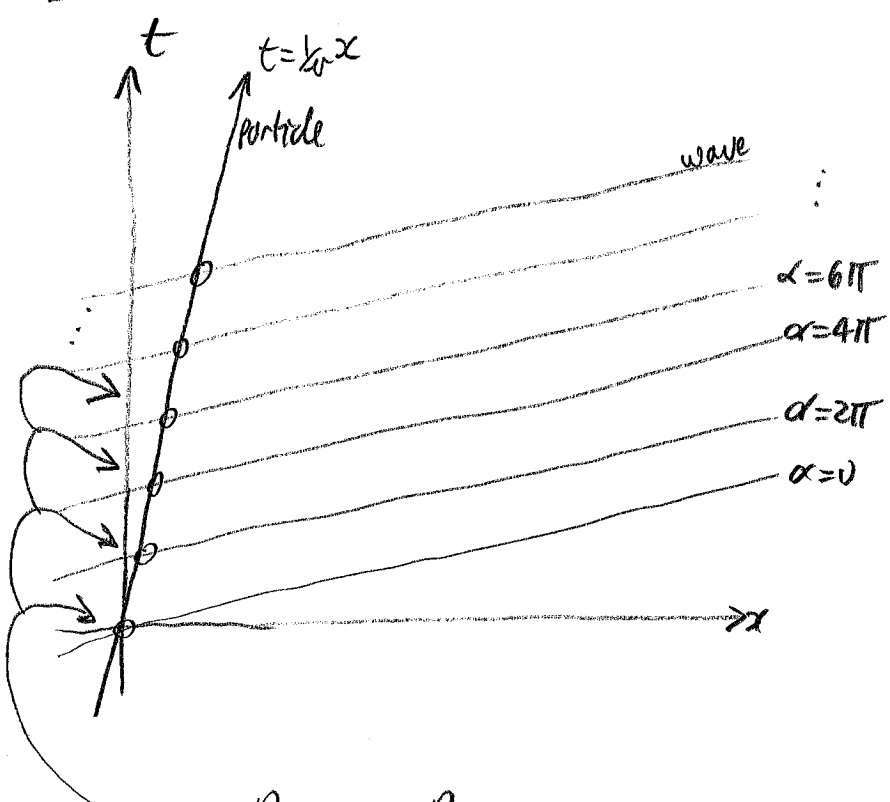
Phase = 0 = $2\pi(\delta\nu_0 t - \frac{\delta\nu_0 v}{c} x)$

$\therefore x = \frac{\delta\nu_0}{\delta\nu_0 \frac{v}{c}} t = \frac{c^2}{v} x$

phase velocity is $v = c^2/\nu$



Superimpose particle and wave



Points of good phase when
 $\alpha = 2\pi n \quad n = 0, 1, 2, \dots$
 i.e.
 $2\pi n = 2\pi v_0 \delta (t - \frac{v_0}{c} x)$

Intersect particle world line when $x = vt$

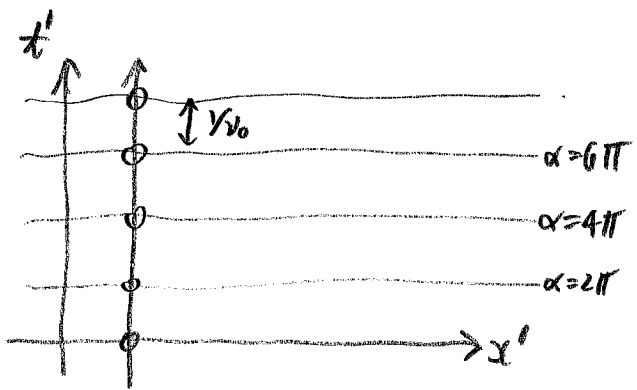
$$n = v_0 \delta (t - \frac{v_0}{c} vt) = v_0 \sqrt{1 - \frac{v^2}{c^2}} t$$

$$t_n = \frac{n}{v_0 \sqrt{1 - \frac{v^2}{c^2}}} = \frac{n}{v_1}$$

i.e. in unprimed frame, flashes separated by time v_1

Easy way to see result

① Particle & wave in particle rest frame



② Transform to unprimed frame

$$\phi = \sin \alpha = \sin(2\pi v_0 t')$$

De Broglie's: Group velocity of the phase waves is equal to the velocity of the moving particle

Add two waves of adjacent frequency $\nu, \nu + \delta\nu$

Use identity

$$\begin{aligned} & \sin(x) + \sin(x + \delta x) \\ &= \sin(\bar{x} - \frac{\delta x}{2}) + \sin(\bar{x} + \frac{\delta x}{2}) \quad \bar{x} = x + \frac{\delta x}{2} \\ &= \sin(\bar{x}) \cos(-\frac{\delta x}{2}) + \cos(\bar{x}) \sin(-\frac{\delta x}{2}) \\ & \quad + \sin(\bar{x}) \cos(+\frac{\delta x}{2}) + \cos(\bar{x}) \sin(+\frac{\delta x}{2}) \end{aligned}$$

cancel $\sin(-\frac{\delta x}{2}) = -\sin(\frac{\delta x}{2})$

$$= 2 \sin \bar{x} \cdot \cos \frac{\delta x}{2} \approx 2 \sin x \cos \frac{\delta x}{2}$$

Apply to waveform

$$\sin 2\pi(\nu t - \frac{\nu}{V} x) \quad \text{to get "beats"}$$

combined form is

$$2 \sin 2\pi(\nu t - \frac{\nu}{V} x) \cdot \cos 2\pi(\delta\nu t - x \delta(\frac{\nu}{V}))$$

sine wave propagating with phase velocity $V = \frac{c^2}{\omega}$

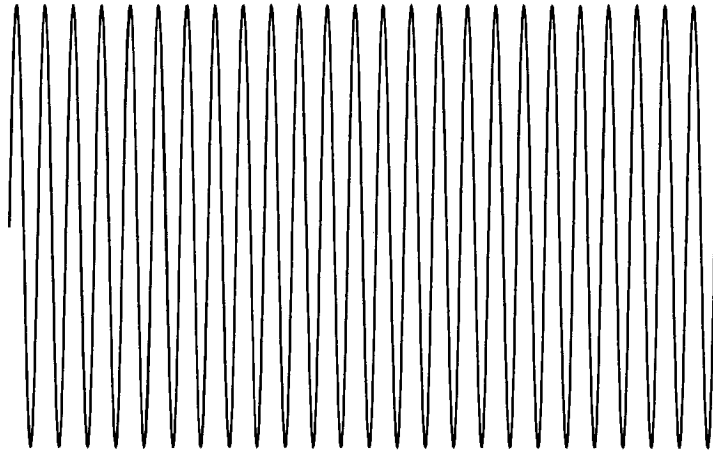
modulating beat envelope propagating at u where

$$\frac{1}{u} = \frac{\delta(\nu/V)}{\delta\nu} \rightarrow \frac{d(\nu/V)}{d\nu}$$

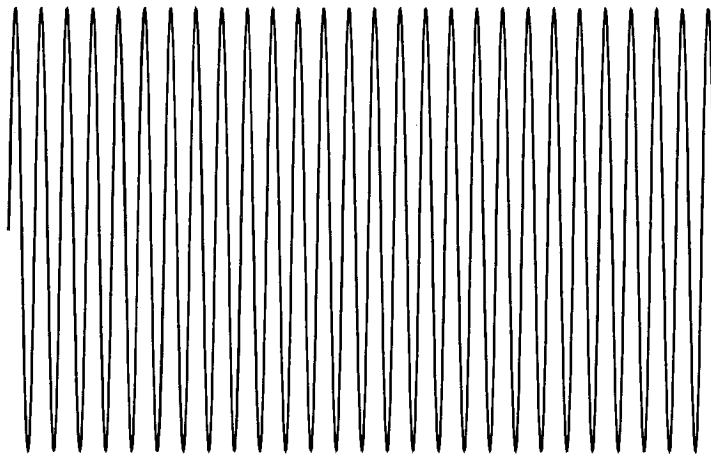
group velocity

"Beats"

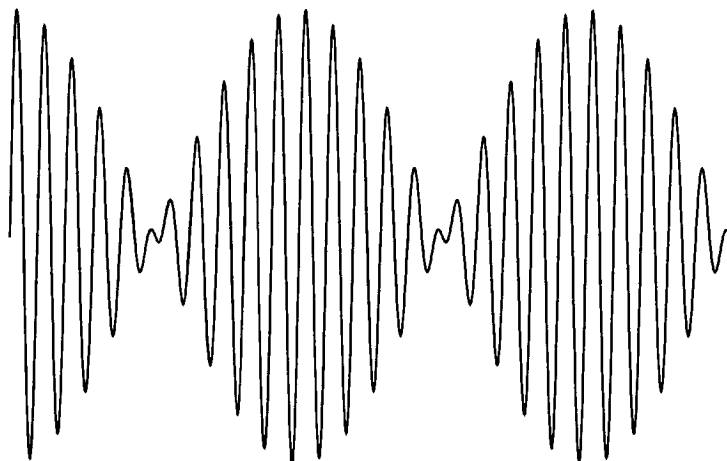
5 1/2



$$\sin(x)$$



$$\sin(1.1x)$$



$$\sin(x) + \sin(1.1x)$$

Compute group velocity

$$\frac{1}{u} = \frac{d(v/\lambda)}{dv}$$

Note functional dependencies!

$$\lambda = \frac{1}{\nu} \frac{m_0 c^2}{\sqrt{1-v^2/c^2}} = \frac{\lambda_0}{\sqrt{1-v^2/c^2}} \quad \nu = \frac{c^2}{v} \quad \frac{1}{\nu} = \frac{v}{c^2}$$

Reconfigure so v is the independent variable:

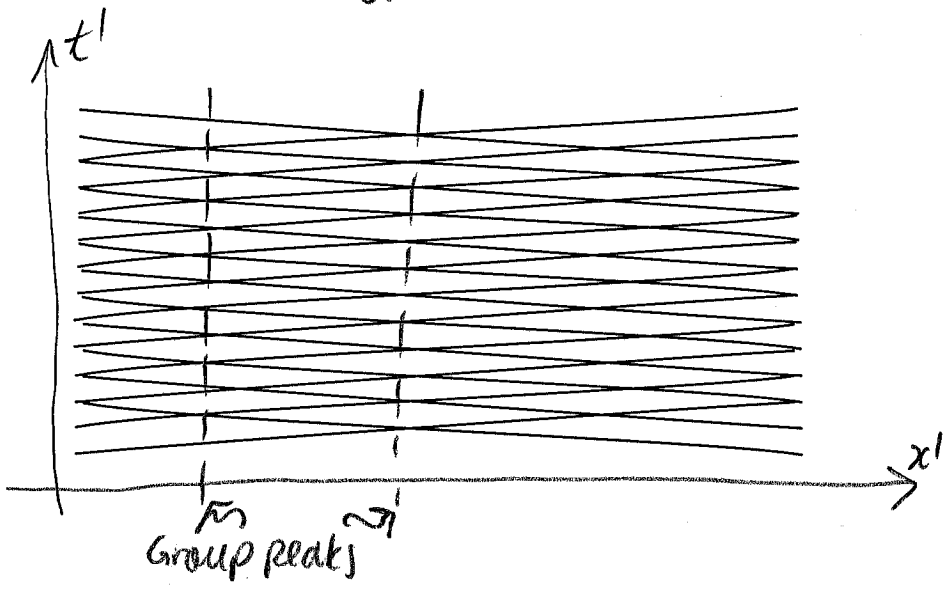
$$\left[\begin{aligned} u &= \frac{dv}{d(v/\lambda)} = \frac{dv/dv}{\frac{d(v/\lambda)}{dv}} = \frac{\lambda_0 \frac{d}{dv} \left(\frac{1}{\sqrt{1-v^2/c^2}} \right)}{\frac{\nu_0}{c^2} \frac{d}{dv} \left(\frac{v}{\sqrt{1-v^2/c^2}} \right)} \quad \textcircled{A} \\ &= \frac{\lambda_0 \frac{v/c^2}{(1-v^2/c^2)^{3/2}}}{\frac{\nu_0}{c^2} \frac{1}{(1-v^2/c^2)^{3/2}}} = v \quad \textcircled{B} \end{aligned} \right] \quad \text{***}$$

$$\text{since } \textcircled{A} \quad \frac{d}{dv} \frac{1}{(1-v^2/c^2)^{1/2}} = \frac{-\frac{1}{2}}{(1-v^2/c^2)^{3/2}} \cdot \frac{-2v}{c^2} = \frac{v/c^2}{(1-v^2/c^2)^{3/2}}$$

$$\begin{aligned} \textcircled{B} \quad \frac{d}{dv} \frac{v}{(1-v^2/c^2)^{1/2}} &= \frac{1}{(1-v^2/c^2)^{1/2}} + \frac{v^2/c^2}{(1-v^2/c^2)^{3/2}} = \frac{(1-v^2/c^2)}{(1-v^2/c^2)^{3/2}} + \frac{v^2/c^2}{(1-v^2/c^2)^{3/2}} \\ &= \frac{1}{(1-v^2/c^2)^{3/2}} \end{aligned}$$

Group velocity / Phase velocities relationship in a picture

Add two plane waves: one moves at $\frac{c^2}{sv} \approx +\infty$ one moves at $\frac{c^2}{-sv} \approx -\infty$ sv very small



By symmetry, interference pattern produces a wave at rest.

Transform to frame in which peaks move at v

