

# HW #4 Chem 1410

1. Show that  $\hat{H}$  for the hydrogen atom commutes with  $\hat{L}^2$  and  $\hat{L}_z$ .

$$\hat{H} = -\frac{\hbar^2}{2m_e r^2} \frac{\partial}{\partial r} \left( r^2 \frac{\partial}{\partial r} \right) - \frac{e^2}{4\pi\epsilon_0 r} + \frac{\hat{L}^2}{2m_e r^2}$$

$[\hat{L}^2, \hat{H}] = 0$  since  $\hat{L}^2$  is independent of  $r$  and  $\hat{L}^2$  commutes with itself.

$[\hat{L}_z, \hat{H}] = \frac{1}{2m_e r^2} [\hat{L}_z, \hat{L}^2]$  since  $\hat{L}_z$  is independent of  $r$

$\hat{L}_z = \frac{\hbar}{i} \frac{\partial}{\partial \phi}$  commutes with  $\hat{L}^2$  because it does nothing to the terms that depend on  $\theta$  and  $[\frac{\partial}{\partial \phi}, \frac{\partial^2}{\partial \phi^2}] = 0$

$\Rightarrow [\hat{L}_z, \hat{H}] = 0$

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2. Is the  $2p_x$  orbital of the H atom an eigenfunction of  $\hat{L}_z$ ? Of  $\hat{L}^2$ ?

We need only consider the angular part of the orbital.  $2p_x \sim \sin\theta \cos\phi$

$\hat{L}_z = \frac{\hbar}{i} \frac{\partial}{\partial \phi} \Rightarrow \hat{L}_z(2p_x) \sim -\sin\theta \sin\phi$

This is not an eigenfunction of  $\hat{L}_z$

$$\begin{aligned}
\hat{L}^2 \sin\theta \cos\phi &= -\hbar^2 \left\{ \frac{1}{\sin\theta} \frac{\partial}{\partial\theta} \left( \sin\theta \frac{\partial}{\partial\theta} \right) + \frac{1}{\sin^2\theta} \frac{\partial^2}{\partial\phi^2} \right\} \sin\theta \cos\phi \\
&= -\hbar^2 \left\{ \frac{\cos\phi}{\sin\theta} \frac{\partial}{\partial\theta} (\sin\theta \cos\theta) + \frac{1}{\sin^2\theta} \sin\theta (-\cos\phi) \right\} \\
&= -\hbar^2 \left\{ \frac{\cos\phi}{\sin\theta} (\cos^2\theta - \sin^2\theta) + \frac{-1}{\sin\theta} \cos\phi \right\} \\
&= -\hbar^2 \left\{ \frac{\cos^2\theta - \sin^2\theta - 1}{\sin\theta} \right\} \cos\phi = -\hbar^2 \left( \frac{-2\sin^2\theta}{\sin\theta} \right) \cos\phi \\
&= 2\hbar^2 \sin\theta \cos\phi
\end{aligned}$$

So the  $2p_x$  orbital is an eigenfunction of  $\hat{L}^2$ .

Now consider the  $2p_+$  orbital: Is it an eigenfunction of  $\hat{L}_z$ ? of  $\hat{L}^2$ ?

$$\hat{L}_z \sin\theta e^{i\phi} = \frac{\hbar}{i} \sin\theta (i) e^{i\phi} = \hbar \sin\theta e^{i\phi}$$

Yes,  $2p_+$  is an eigenfunction of  $\hat{L}_z$ .

$$\begin{aligned}
\hat{L}^2 \sin\theta e^{i\phi} &= -\hbar^2 \left\{ \frac{\cos^2\theta - \sin^2\theta}{\sin\theta} + \frac{1}{\sin\theta} (i)^2 \right\} e^{i\phi} \\
&= -\hbar^2 \left\{ \frac{\cos^2\theta - \sin^2\theta - 1}{\sin\theta} \right\} e^{i\phi} \\
&= 2\hbar^2 \sin\theta e^{i\phi}
\end{aligned}$$

Yes,  $2p_+$  is an eigenfunction of  $\hat{L}^2$ .

3. Show that the  $1s \rightarrow 2s$  transition of the H atom is dipole forbidden.

The operator responsible for a dipole transition is proportional to  $x, y,$  or  $z$ .

$$x = r \sin \theta \cos \phi$$

$$y = r \sin \theta \sin \phi$$

$$z = r \cos \theta$$

The integrals factor into  $r, \theta, \phi$  integrals. We can show that they are zero by examining only the  $\theta$  or  $\phi$  parts.

$$\begin{aligned} \langle 1s | z | 2s \rangle &\propto \int_0^\pi \sin \theta \cos \theta d\theta \int_0^{2\pi} d\phi \\ &= 2\pi \int_0^\pi \cos \theta d(\cos \theta) = \frac{2\pi \cos^2 \theta}{2} \Big|_0^\pi = 0 \end{aligned}$$

$$\begin{aligned} \langle 1s | x | 2s \rangle &\propto \int_0^{2\pi} \cos \phi d\phi \int_0^\pi \sin^2 \theta d\theta \\ \int_0^{2\pi} \cos \phi d\phi &= \sin \phi \Big|_0^{2\pi} = 0 \end{aligned}$$

$$\text{So } \langle 1s | x | 2s \rangle = 0$$

$$\begin{aligned} \langle 1s | y | 2s \rangle &\propto \int_0^{2\pi} \sin \phi d\phi \int_0^\pi \sin^2 \theta d\theta \\ \int_0^{2\pi} \sin \phi d\phi &= -\cos \phi \Big|_0^{2\pi} = 0 \end{aligned}$$

$$\text{So } \langle 1s | y | 2s \rangle = 0$$

4. What is the average of  $r$  for the  $2s$  orbital of the  $H$  atom? How does this differ from the most probable value of  $r$ ?

$$\begin{aligned}
 \langle r \rangle_{2s} &= \int_0^{\infty} dr \int_0^{\pi} \sin\theta d\theta \int_0^{2\pi} d\phi r^3 \psi_{2s}^2 \\
 &= \frac{4\pi}{32\pi a_0^3} \int_0^{\infty} r^3 \left(2 - \frac{r}{a_0}\right)^2 e^{-r/a_0} dr \\
 &= \frac{1}{8a_0^3} \int_0^{\infty} r^3 \left(4 - \frac{4r}{a_0} + \frac{r^2}{a_0^2}\right) e^{-r/a_0} dr \\
 &= \frac{1}{8a_0^3} \left\{ \int_0^{\infty} 4r^3 e^{-r/a_0} dr - \frac{4}{a_0} \int_0^{\infty} r^4 e^{-r/a_0} dr + \frac{1}{a_0^2} \int_0^{\infty} r^5 e^{-r/a_0} dr \right\} \\
 &= \frac{1}{8a_0^3} \left\{ 4 \cdot 3! a_0^4 - \frac{4}{a_0} \cdot 4! a_0^5 + \frac{1}{a_0^2} \cdot 5! a_0^6 \right\} \\
 &= \frac{a_0}{8} \{ 24 - 96 + 120 \} = 6a_0
 \end{aligned}$$

Now let's find the most probable value of  $r$ .

$$P(r) \propto r^2 \psi^2 = r^2 \left(4 - \frac{4r}{a_0} + \frac{r^2}{a_0^2}\right) e^{-r/a_0}$$

$$P(r) \propto \left(4r^2 - \frac{4r^3}{a_0} + \frac{r^4}{a_0^2}\right) e^{-r/a_0}$$

$$\frac{dP}{dr} = 0 = \left(8r - \frac{12r^2}{a_0} + \frac{4r^3}{a_0^2}\right) e^{-r/a_0} - \left(\frac{1}{a_0}\right) \left(4r^2 - \frac{4r^3}{a_0} + \frac{r^4}{a_0^2}\right) e^{-r/a_0}$$

$$\begin{aligned}
 0 &= 8r - \frac{12r^2}{a_0} + \frac{4r^3}{a_0^2} \\
 &\quad - \frac{4r^2}{a_0} + \frac{4r^3}{a_0^2} - \frac{r^4}{a_0^3}
 \end{aligned}$$

$$8r - \frac{16r^2}{a_0} + \frac{8r^3}{a_0^2} - \frac{r^4}{a_0^3} = 0$$

$$\Rightarrow r \left( 8 - \frac{16r}{a_0} + \frac{8r^2}{a_0^2} - \frac{r^3}{a_0^3} \right) = 0$$

$$\Rightarrow r = 0 \leftarrow \text{one solution}$$

$$\frac{r^3}{a_0^3} - \frac{8r^2}{a_0^2} + \frac{16r}{a_0} - 8 = 0 \quad \left. \vphantom{\frac{r^3}{a_0^3}} \right\} \text{let } x = \frac{r}{a_0}$$

$$x^3 - 8x^2 + 16x - 8$$

By inspection,  $x=2$  is a solution.

If we factor this out, get

$$x^2 - 6x + 4 = 0 \Rightarrow x = 3 \pm \frac{1}{2}\sqrt{20} = 0.77, 5.23$$

So the extrema occur at

$$r_1 = 0$$

$$r_2 = 0.77a_0$$

$$r_3 = 2a_0$$

$$r_4 = 5.23a_0$$

