

Chapter 2

Consider a multiphase, multi-component system

$$E = \sum_{\alpha=1}^{\nu} E^{(\alpha)}, \quad \nu = \# \text{ phases}$$

Different phases = different subsystems

$$S = \sum_{\alpha=1}^{\nu} S^{(\alpha)}$$

$$V = \sum_{\alpha=1}^{\nu} V^{(\alpha)}$$

$$n_i = \sum_{\alpha=1}^{\nu} n_i^{(\alpha)}, \quad n_i^{(\alpha)} = \# \text{ of moles of species } i \text{ in phase } \alpha$$

$$\delta E = \sum_{\alpha=1}^{\nu} T^{(\alpha)} \delta S^{(\alpha)} - p^{(\alpha)} \delta V^{(\alpha)} + \sum_{i=1}^r \mu_i^{(\alpha)} \delta n_i^{(\alpha)} \quad \left| \quad r = \# \text{ species} \right.$$

$$\begin{array}{l} \text{At equilibrium} \\ (\delta E)_{S,V,n_i} \geq 0 \end{array} \quad \left| \quad \begin{array}{l} \sum \delta S^{(\alpha)} = 0 \\ \sum \delta V^{(\alpha)} = 0 \\ \sum \delta n_i^{(\alpha)} = 0, \quad i = 1, 2, \dots, r \end{array} \right.$$

This follows from the fact the total S, V, n are constant

Special case $\nu = 2$

$$\delta S^{(1)} = -\delta S^{(2)}$$

$$\delta V^{(1)} = -\delta V^{(2)}$$

$$\delta n_j^{(1)} = -\delta n_j^{(2)}$$

$$\delta E = (T^{(1)} - T^{(2)})\delta S^{(1)} - (p^{(1)} - p^{(2)})\delta V^{(1)} + \sum_{i=1}^r (\mu_i^{(1)} - \mu_i^{(2)})\delta n_i^{(1)}$$

Must hold for all small **unconstrained** variations $\delta S^{(1)}$, $\delta V^{(1)}$, $\delta n_i^{(1)}$

and since $\delta E \geq 0$

$$\Rightarrow T^{(1)} = T^{(2)}, p^{(1)} = p^{(2)}, \mu_i^{(1)} = \mu_i^{(2)}$$

which implies $(\delta E)_{S,V,n_i} = 0$

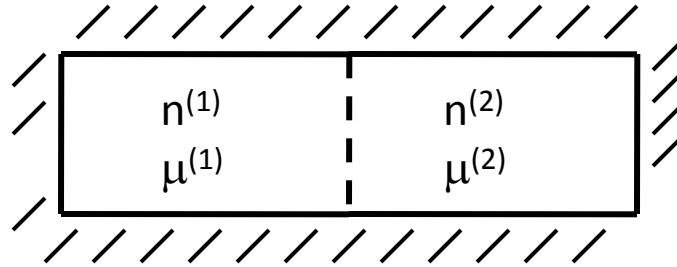
for small displacements away from equilibrium

Mass equilibrium: $\mu^{(1)} = \mu^{(2)}$

If $\mu^{(1)} > \mu^{(2)}$

mass will flow until
equil. is achieved

$$\mu_{fin}^{(1)} = \mu_{fin}^{(2)}$$



$$\Delta S > 0$$

Assume no work done
and no heat flow

$$\Delta S = -\frac{\mu^{(1)}}{T} \Delta n^{(1)} - \frac{\mu^{(2)}}{T} \Delta n^{(2)} = -\frac{\mu^{(1)} - \mu^{(2)}}{T} \Delta n^{(1)}$$

$$\Delta S > 0 \Rightarrow \Delta n^{(1)} < 0$$

matter flows from high μ to low μ

gradients in $\frac{\mu}{T} \rightarrow$ mass flow

gradients in $\frac{1}{T} \rightarrow$ force causing
heat flow

gradients in $\frac{P}{T} \rightarrow$ volume change

$(\Delta E)_{S,V,n} > 0$ for displacements away from equilibrium
if unconstrained $(\delta E)_{S,V,n} = 0$

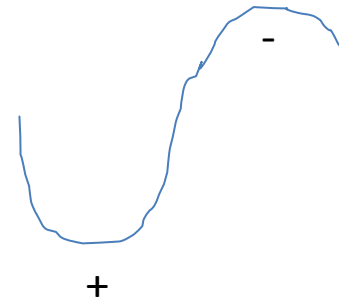
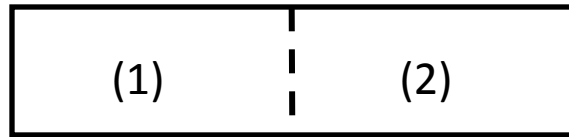
$(\Delta E)_{S,V,n} = (\delta^2 E)_{S,V,n} + (\delta^3 E)_{S,V,n} + \dots$ near equilibrium

$$(\delta^{(2)} E)_{S,V,n} \geq 0$$

$(\delta^{(2)} E)_{S,V,n} > 0 \Rightarrow$ Stable wrt small fluctuations
away from equilibrium

$(\delta^{(2)} E)_{S,V,n} = 0 \Rightarrow$ need to examine higher order variations

$(\delta^{(2)}E)_{S,V,n} < 0 \Rightarrow$ system is unstable wrt fluctuations



Consider the situation $\left\{ \begin{array}{l} \delta S = 0 = \delta S^{(1)} + \delta S^{(2)} \\ \delta V^{(1)} = \delta V^{(2)} = 0 \\ \delta n^{(1)} = \delta n^{(2)} = 0 \end{array} \right.$

$$\delta^2 E = (\delta^2 E)^{(1)} + (\delta^2 E)^{(2)} = \frac{1}{2} \left(\frac{\partial^2 E}{\partial S^2} \right)_{V,n}^{(1)} (\delta S^{(1)})^2 + \frac{1}{2} \left(\frac{\partial^2 E}{\partial S^2} \right)_{V,n}^{(2)} (\delta S^{(2)})^2$$

$$\delta^2 E = \frac{1}{2} \left\{ \left(\frac{\partial^2 E}{\partial S^2} \right)_{V,n}^{(1)} + \left(\frac{\partial^2 E}{\partial S^2} \right)_{V,n}^{(2)} \right\} [\delta S^{(1)}]^2$$

$$\left(\frac{\partial^2 E}{\partial S^2}\right)_{V,n} = \left(\frac{\partial T}{\partial S}\right)_{V,n} = \frac{T}{C_V}$$

$$\begin{aligned}(\delta^2 E)_{S,V,n} &= \frac{1}{2} \left\{ \frac{T^{(1)}}{C_V^{(1)}} + \frac{T^{(2)}}{C_V^{(2)}} \right\} (\delta S^{(1)})^2 \\ &= \frac{(\delta S^{(1)})^2}{2} T \left[\frac{1}{C_V^{(1)}} + \frac{1}{C_V^{(2)}} \right]\end{aligned}$$

$$(\delta^{(2)} E)_{S,V,n} \geq 0 \Rightarrow T \left[\frac{1}{C_V^{(1)}} + \frac{1}{C_V^{(2)}} \right] \geq 0$$

this should hold for any subdivision

$$\Rightarrow C_V \geq 0 \quad \Bigg| \quad \text{this is essential if the system is stable}$$