

Pecora-Carroll Master stability equation.

$$\dot{z} = A(t)z + (\alpha + i\beta)Kz$$

Find the regions in (α, β) where $z(t) \rightarrow 0$ as $t \rightarrow \infty$

If we write $z = R + iS$ then

$$\frac{dR}{dt} = A(t)R + \alpha KR - \beta KS$$

$$\frac{dS}{dt} = A(t)S + \alpha KS + \beta KR$$

Can just integrate this or ~~compute~~ for periodic $A(t)$ you can compute the monodromy matrix as a function of (α, β)

Surprisingly, even for oscillators it is possible that some kind of coupling lead to instability

For homework, I will give you the following

$$\dot{z}_j = z_j \left(1 - \underbrace{(1+iq)}_{\text{real}} z_j \bar{z}_j \right) + (\beta + i\gamma)(z_k - z_j)$$

q, γ are parameters

In the mean time, consider the following example

Coupled Brusselators:

$$\dot{X}_j = A - (B+1)X_j + X_j^2 Y_j + D_x (X_n - X_j)$$

$$\dot{Y}_j = B X_j - X_j^2 Y_j + D_y (Y_n - Y_j)$$

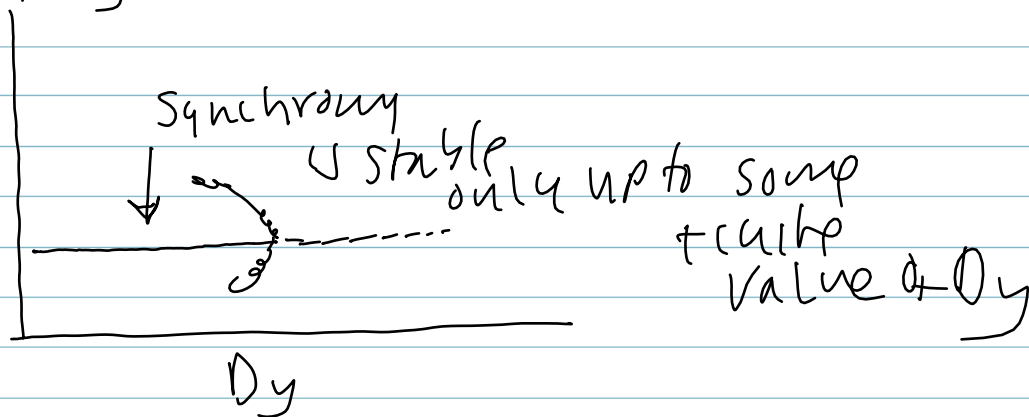
I will use the computer form 11.

Fix $A =$

$B =$

$D_x = 0.05$

Vary D_y



General diffusive coupling

$$\dot{X}_i = F(X_i) + \sum_{j=1}^N \hat{C}_{ij} k(X_j - X_i) \quad \hat{C}_{ii} = 0 \quad i=1, \dots, N \quad (\hat{\star})$$

Where $F: \mathbb{R}^m \rightarrow \mathbb{R}^m$ $k \in \mathbb{R}^{m \times m}$

$$(\hat{C}_{ij}) \equiv \hat{C} \in \mathbb{R}^{N \times N}$$

Let $x_i = u(t)$ be a synchronous solution.

That is

$$\dot{u}(t) = F(u(t)) \quad (S)$$

$$\text{Define } c_{ij} = \begin{cases} \hat{C}_{ij} & i \neq j \\ -\sum_{j=1}^N \hat{C}_{ij} & i = j \end{cases}$$

$$\text{Let } C = (c_{ij}) \quad \text{note } C \begin{bmatrix} 1 \\ \vdots \\ 1 \end{bmatrix} = 0$$

We can rewrite $(\hat{\star})$ as

$$\dot{X}_i = F(X_i) + \sum_{j=1}^N c_{ij} k X_j \quad (\star)$$

$$\text{since } -\sum_{\substack{j=1 \\ j \neq i}}^N \hat{C}_{ij} = c_{ii}$$

To check stability of $x_i = u(t)$ write

$$x_i(t) = u(t) + y_i \Rightarrow (L)$$

$$\dot{y}_i = A(t) y_i + \sum_{j=1}^N c_{ij} k y_j \quad \text{where } A(t) \equiv D_x F(u(t))$$

We can write (1) as

$$\dot{Y} = [I \otimes A(t) + C \otimes K] Y$$

But, let's just stick with (L) for now.

Let $C \vec{v} = \lambda \vec{v}$ be eigenvalue-eigen vector pair for C. Note

$$C \in \mathbb{R}^{N \times N} \quad \& \quad \vec{v} \in \mathbb{R}^N \subset \mathbb{C}^N$$

Write $Y_i(t) = v_i W(t)$ where $(v_1, \dots, v_N)^T \equiv \vec{v}$

Note $W(t) \in \mathbb{C}^m$ for each t.

(Keep this straight!) . Plug this into (L)

to get

$$v_i \dot{W}(t) = v_i [A(t)W(t) + \sum c_{ij} v_j K W(t)]$$

$$= v_i [A(t)W(t) + \lambda K W(t)]$$

Multiply by v_i & sum up to get

$$\|\vec{v}\|^2 \dot{W}(t) = \|\vec{v}\|^2 [A(t)W(t) + \lambda K W(t)]$$

Since $\|\vec{v}\| \neq 0 \Rightarrow$

$$\dot{W}(t) = A(t)W(t) + \lambda K W(t) \quad (\star\star)$$

Thus we've reduced a $N \times m$ dim system to N m -dim systems!

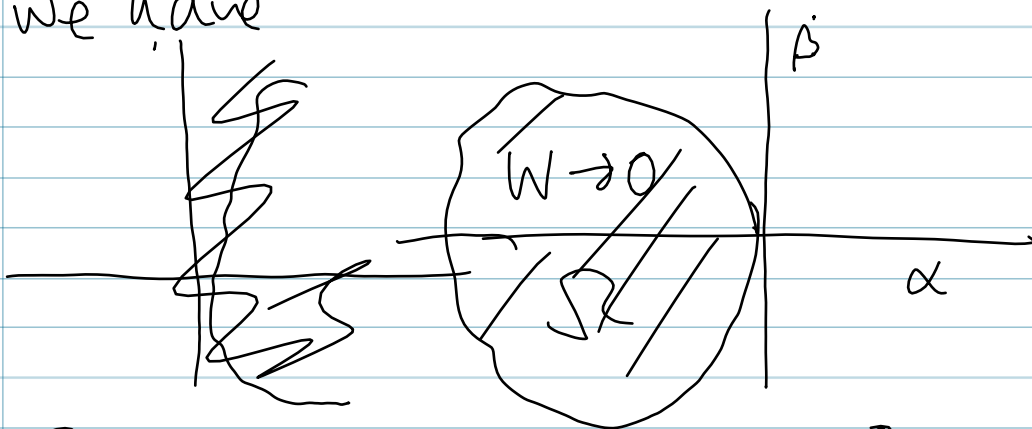
Since $\lambda \in \mathbb{C}$ we need only look at
 $\dot{w}(t) = A(t)w(t) + (\alpha + i\beta)/K w(t)$

as α, β range around complex plane

~~Then for given value~~

choice of

suppose for some set of ~~parameters~~
 we have



Then if $\lambda \in \Omega \Rightarrow w(t) \rightarrow 0$

so we can pick coupling matrix C
 so that all eigenvalues of C are in Ω
 to guarantee that synchrony is stable!

This is the essence of the Perron master
 stability theorem

Unfortunately, you still have to commute!