

# Systematic uncertainties in long baseline neutrino oscillations for large $\theta_{13}$

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Center for Neutrino Physics at Virginia Tech

Based on

P. Coloma, P. Huber, J. Kopp and W. Winter, 1209.5973 [hep-ph]

Flux Measurement and Determination in the Intensity Frontier Era Neutrino Beams

Pittsburgh PACCC, Dec 8, 2012

# Outline

- Introduction to precision in  $\theta_{13}$  and  $\delta$
- Introduction to the problem of systematics for large  $\theta_{13}$ 
  - Possible ways to reduce their impact
- Simulation details
- Results
  - Comparison, key systematics, effect of luminosity

# Precision in $\theta_{13}$

$$(\Delta\theta_{13})_{\pm} \propto \left[ \frac{(1 \mp \hat{A})^2}{\sin^2((1 \mp \hat{A})\Delta)} \right] \frac{1}{\theta_{13}} \Delta N_{\pm}$$

*neutrino/antineutrino*

# Precision in $\theta_{13}$

$$(\Delta\theta_{13})_{\pm} \propto \left[ \frac{(1 \mp \hat{A})^2}{\sin^2((1 \mp \hat{A})\Delta)} \right] \frac{1}{\theta_{13}} \Delta N_{\pm}$$

Statistical limit:

$$\Delta N_{\pm} \propto \sqrt{N_{\pm}} \propto \theta_{13} \longrightarrow (\Delta\theta_{13})_{\pm} \propto \text{const}$$

Systematics on the signal:

$$\Delta N_{\pm} \propto N_{\pm} \propto \theta_{13}^2 \longrightarrow (\Delta\theta_{13})_{\pm} \propto \theta_{13}$$

Background error:

$$\Delta N_{\pm} \propto \text{const} \longrightarrow (\Delta\theta_{13})_{\pm} \propto 1/\theta_{13}$$

# Precision in $\theta_{13}$

Statistical limit:

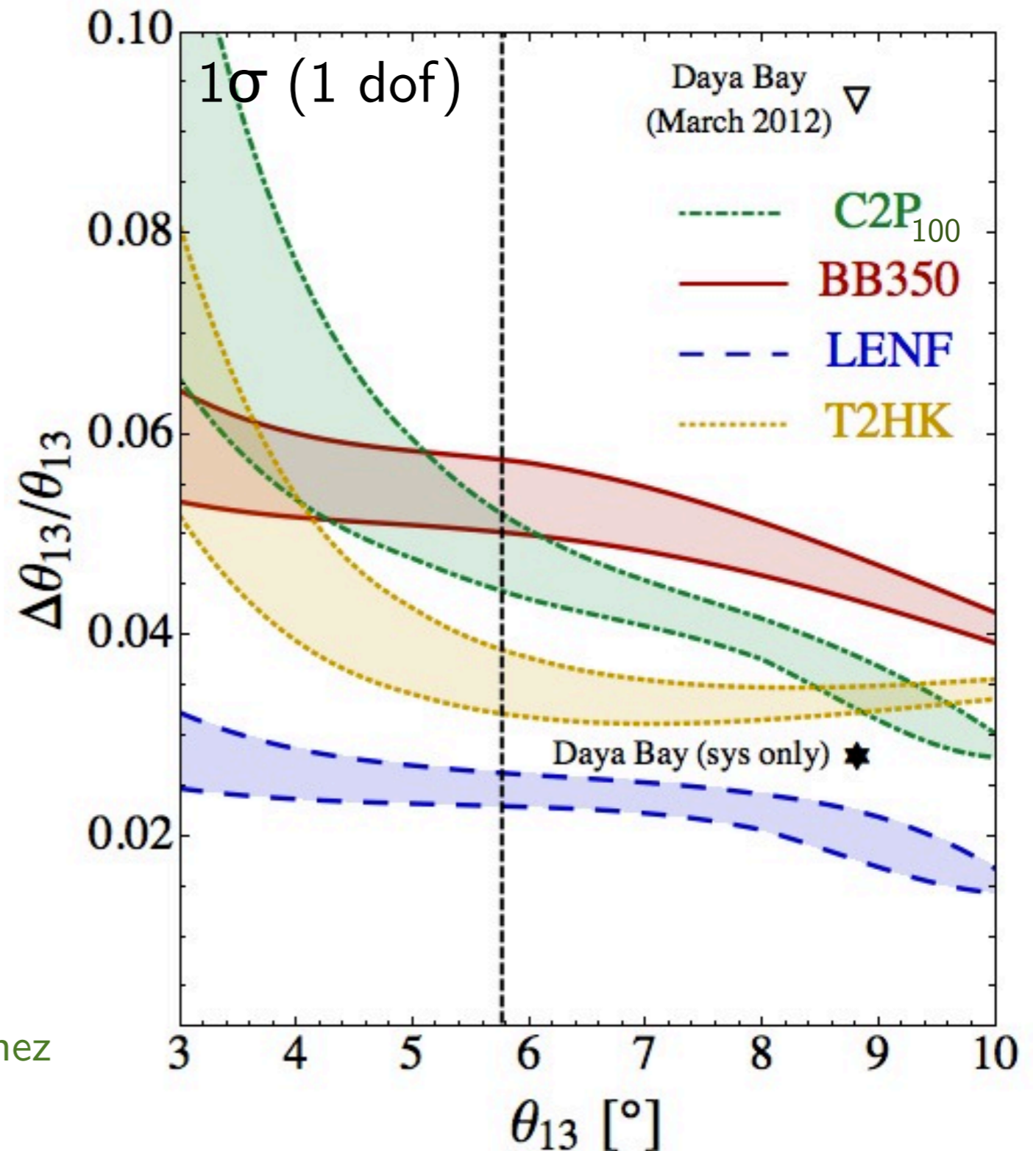
$$\frac{\Delta\theta_{13}}{\theta_{13}} \propto \frac{1}{\theta_{13}}$$

Systematics on the signal:

$$\frac{\Delta\theta_{13}}{\theta_{13}} \propto \text{const}$$

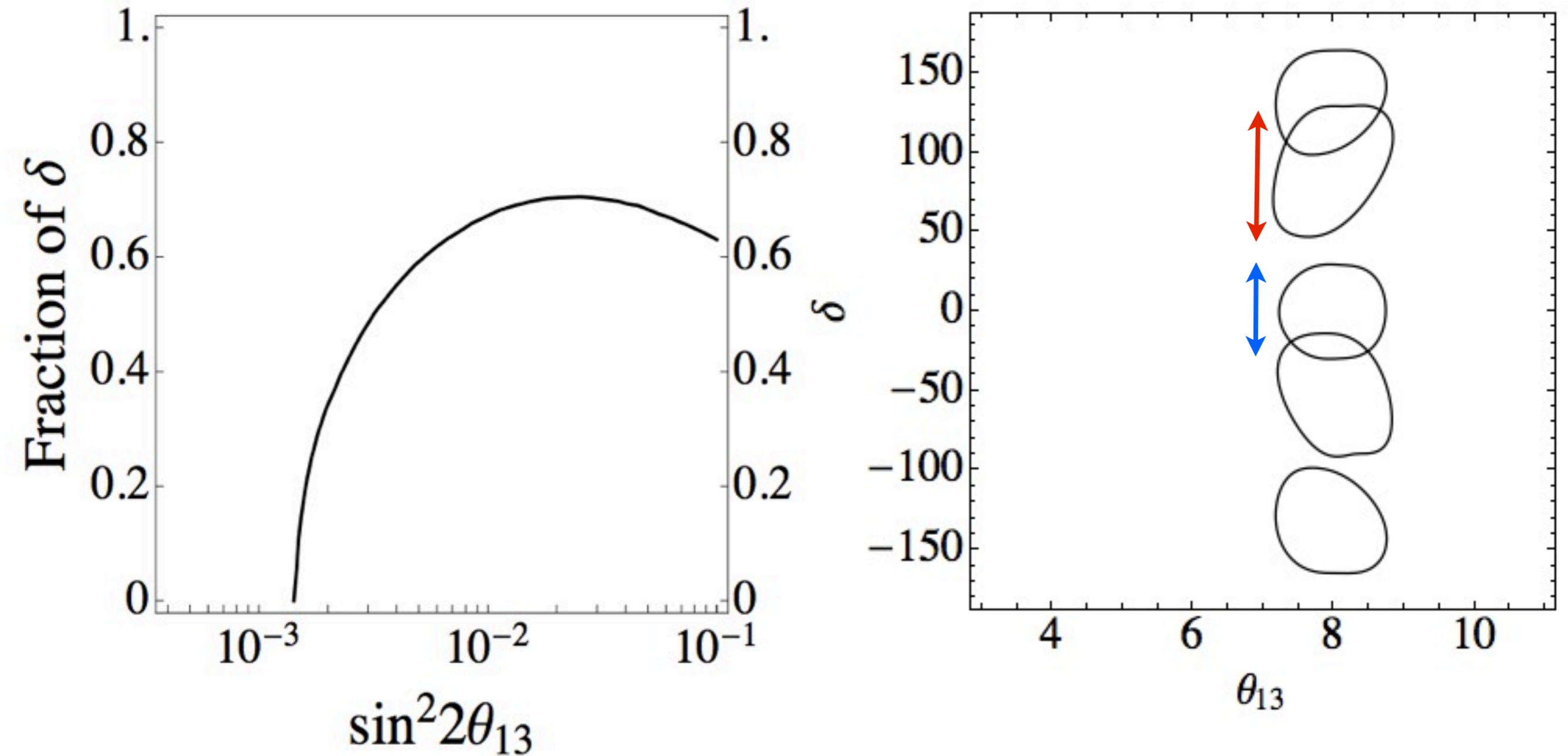
Background error:

$$\frac{\Delta\theta_{13}}{\theta_{13}} \propto \frac{1}{\theta_{13}^2}$$



# Precision in $\delta$

CPV discovery potential vs precision:

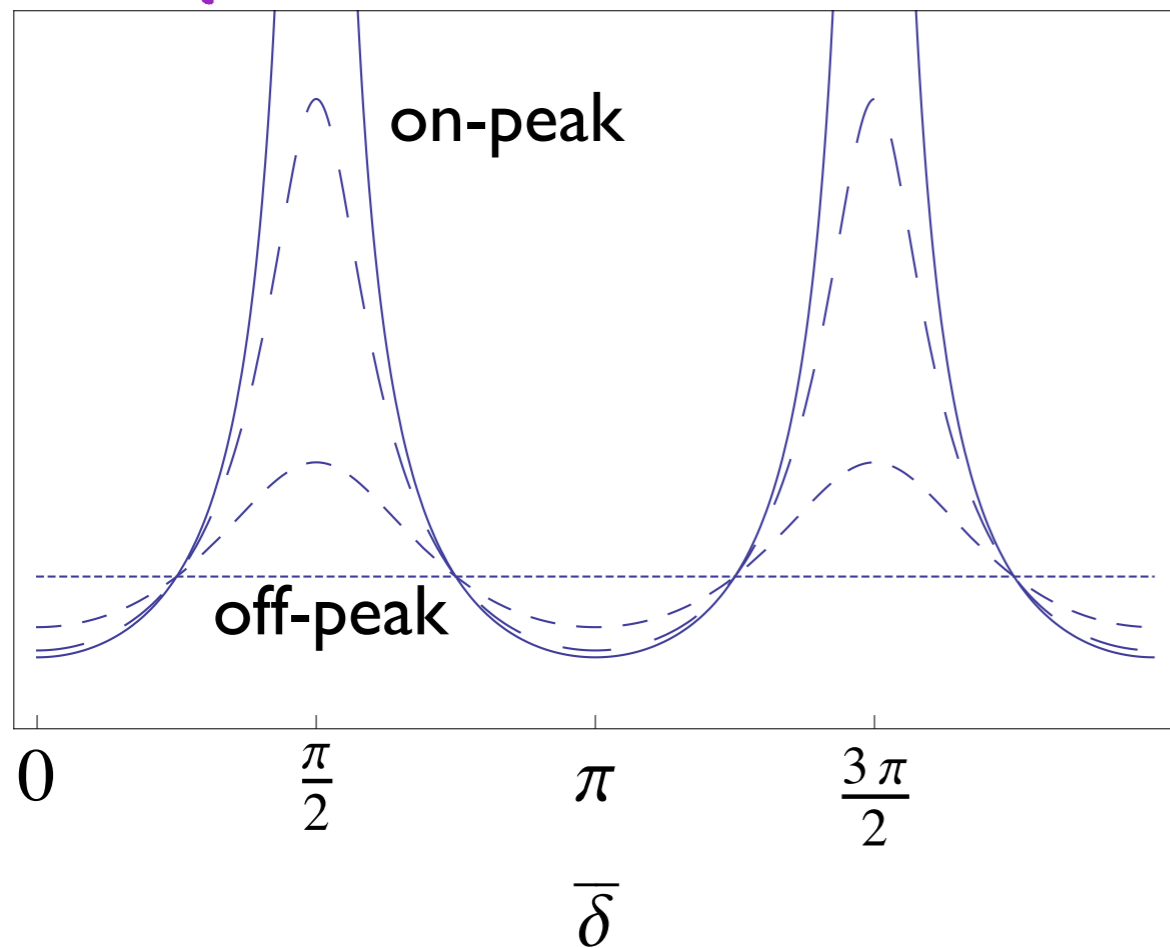


# Precision in $\delta$

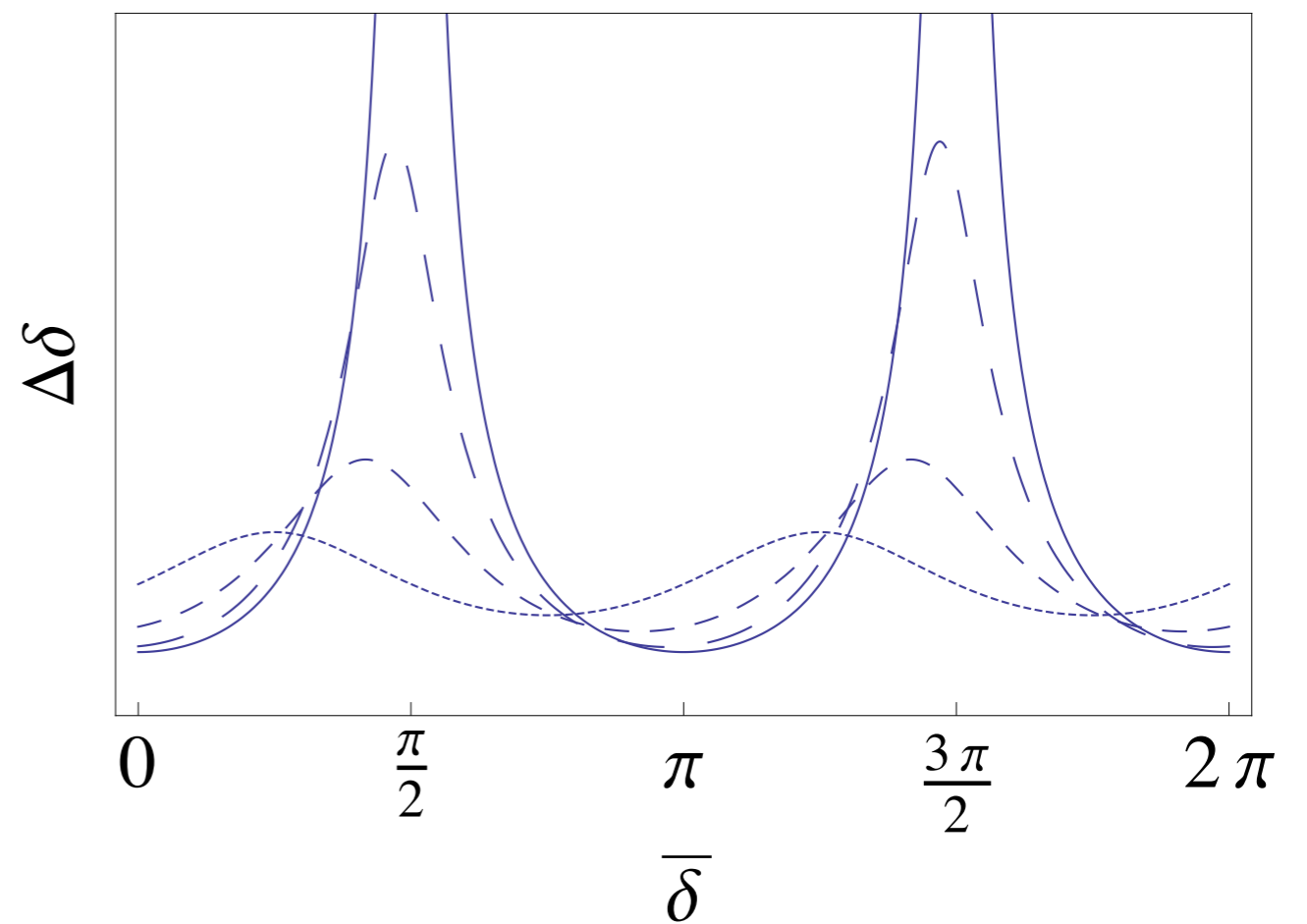
VACUUM

$$(\Delta\delta)_{\pm} \propto f[\Delta] \frac{1}{\sin(\Delta \mp \delta)}$$

*equal*  $\nu$  and  $\bar{\nu}$  events



*less*  $\bar{\nu}$  than  $\nu$  events



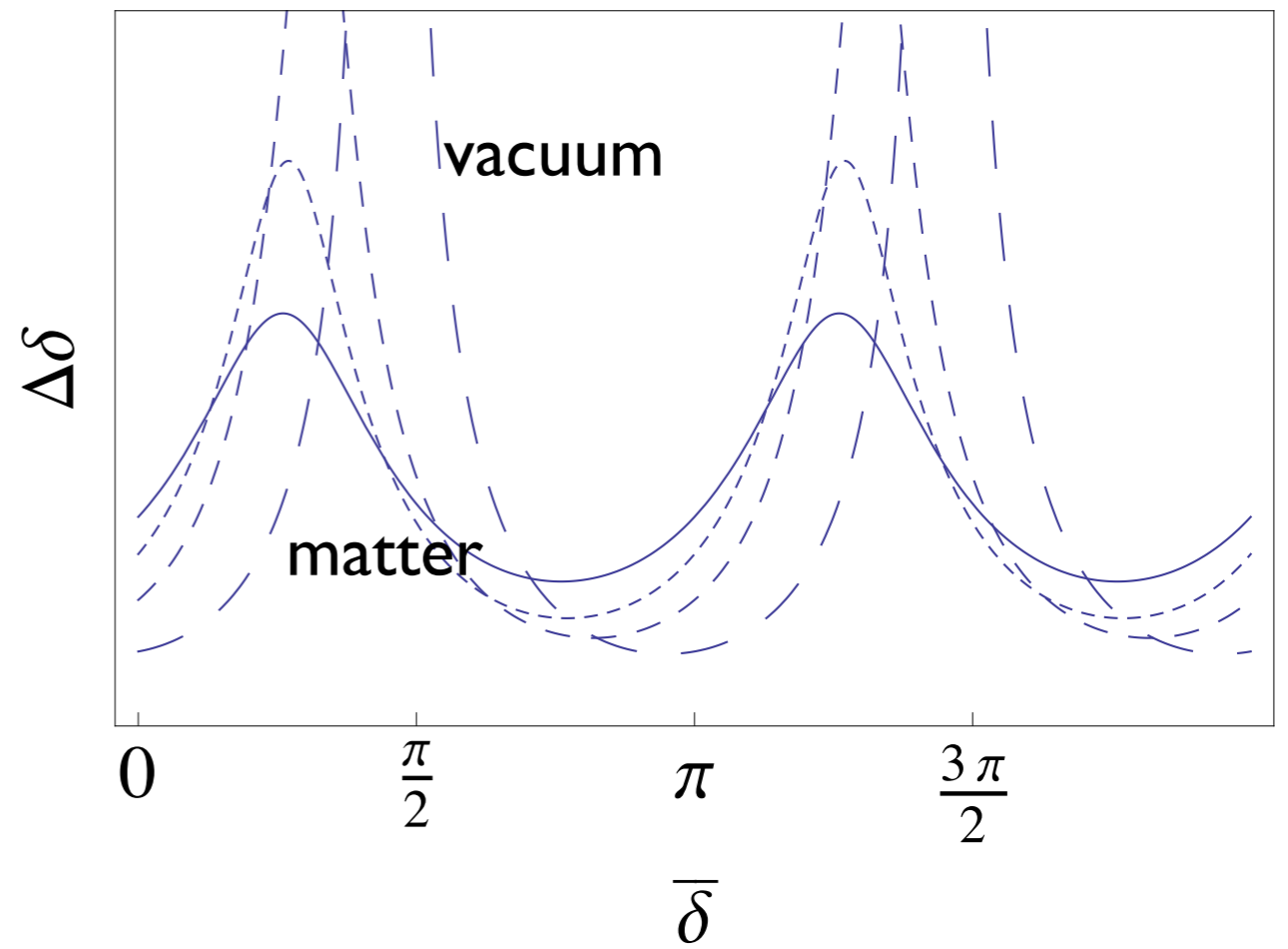
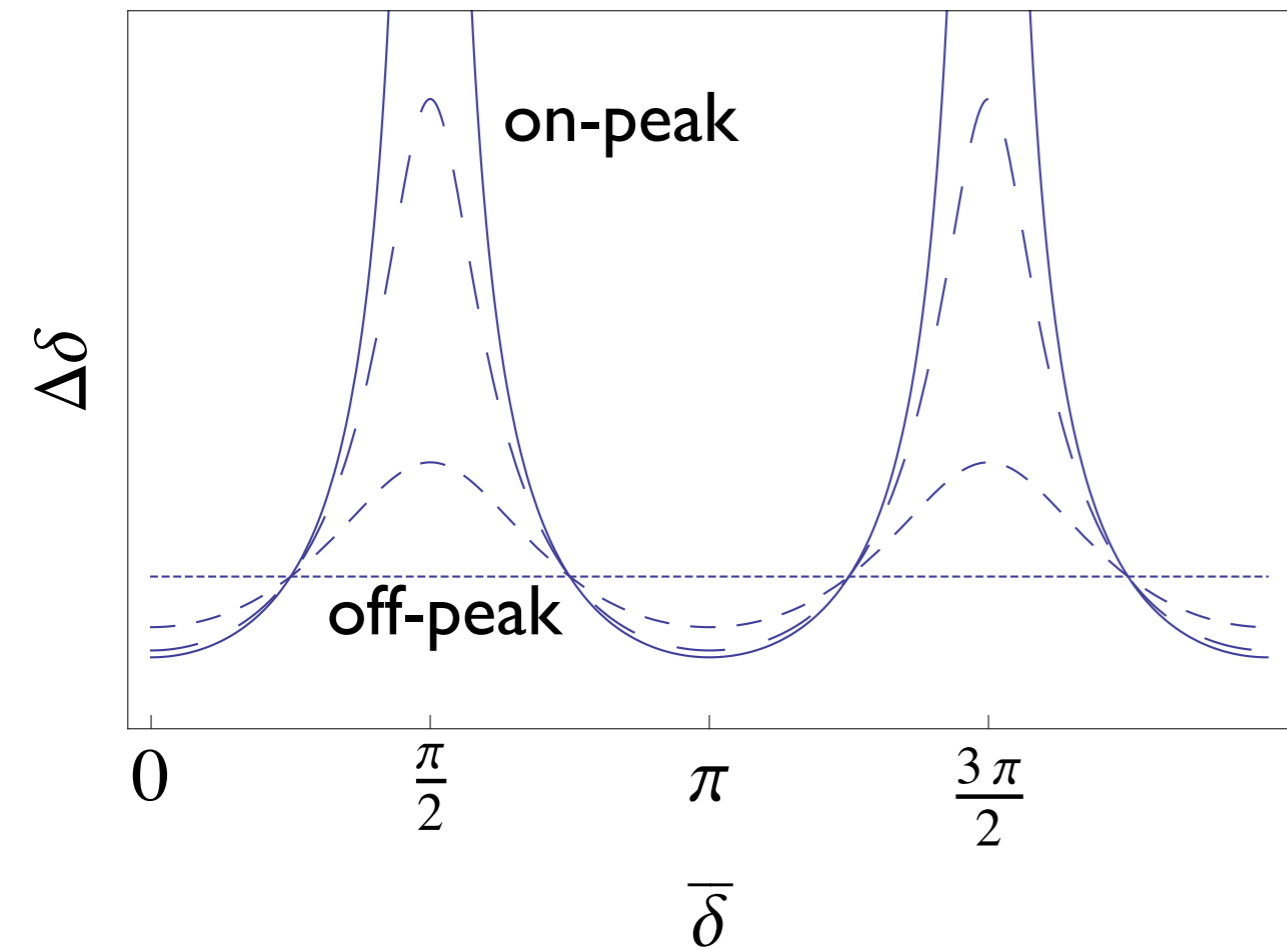
# Precision in $\delta$

VACUUM

$$(\Delta\delta)_{\pm} \propto f[\Delta] \frac{1}{\sin(\Delta \mp \delta)}$$

MATTER

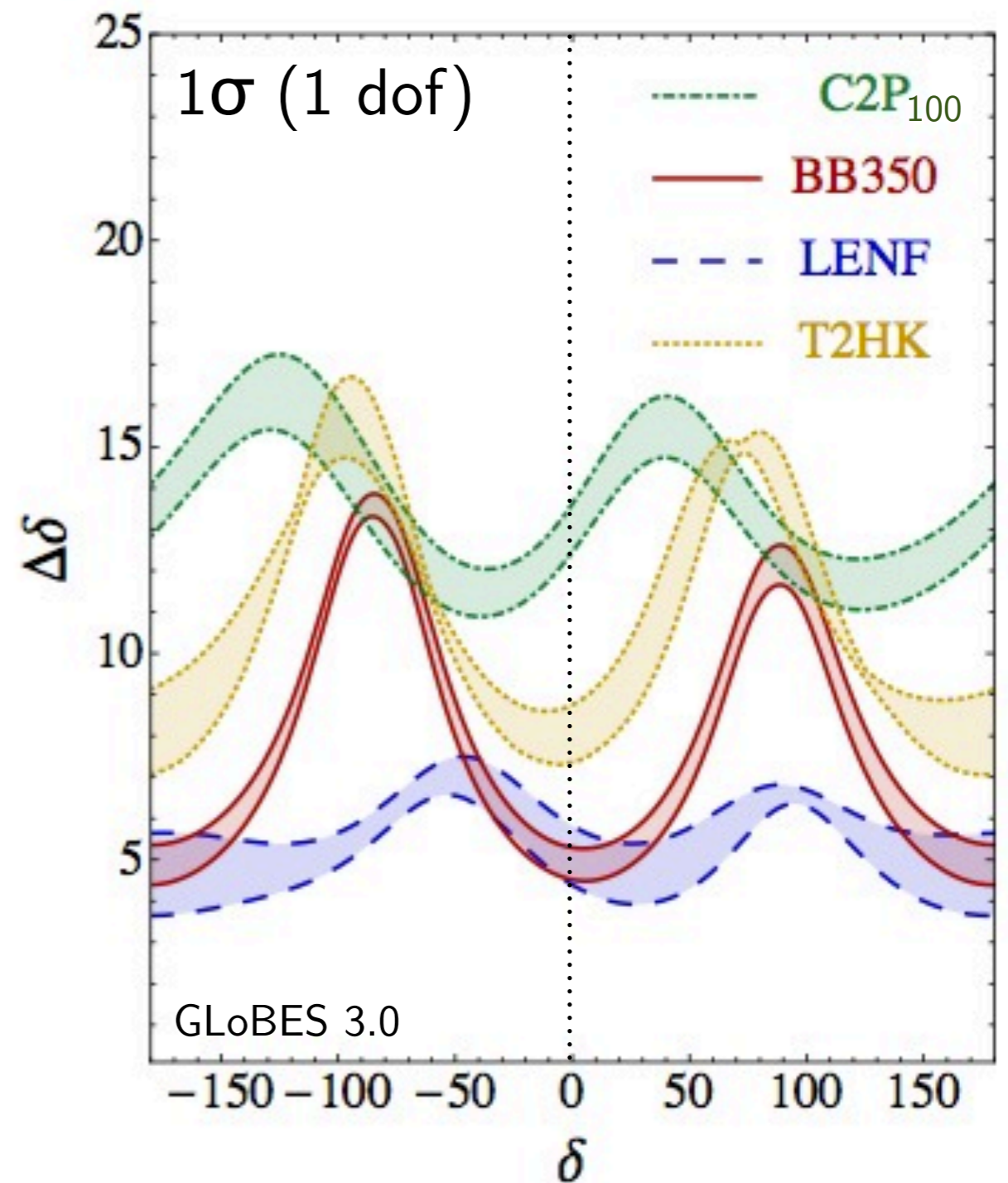
$$(\Delta\delta)_{\pm} \propto \tilde{f}[\hat{A}, \Delta] \frac{1}{\sin\left(\Delta \frac{\hat{A}}{1 \mp \hat{A}} \mp \delta\right)}$$





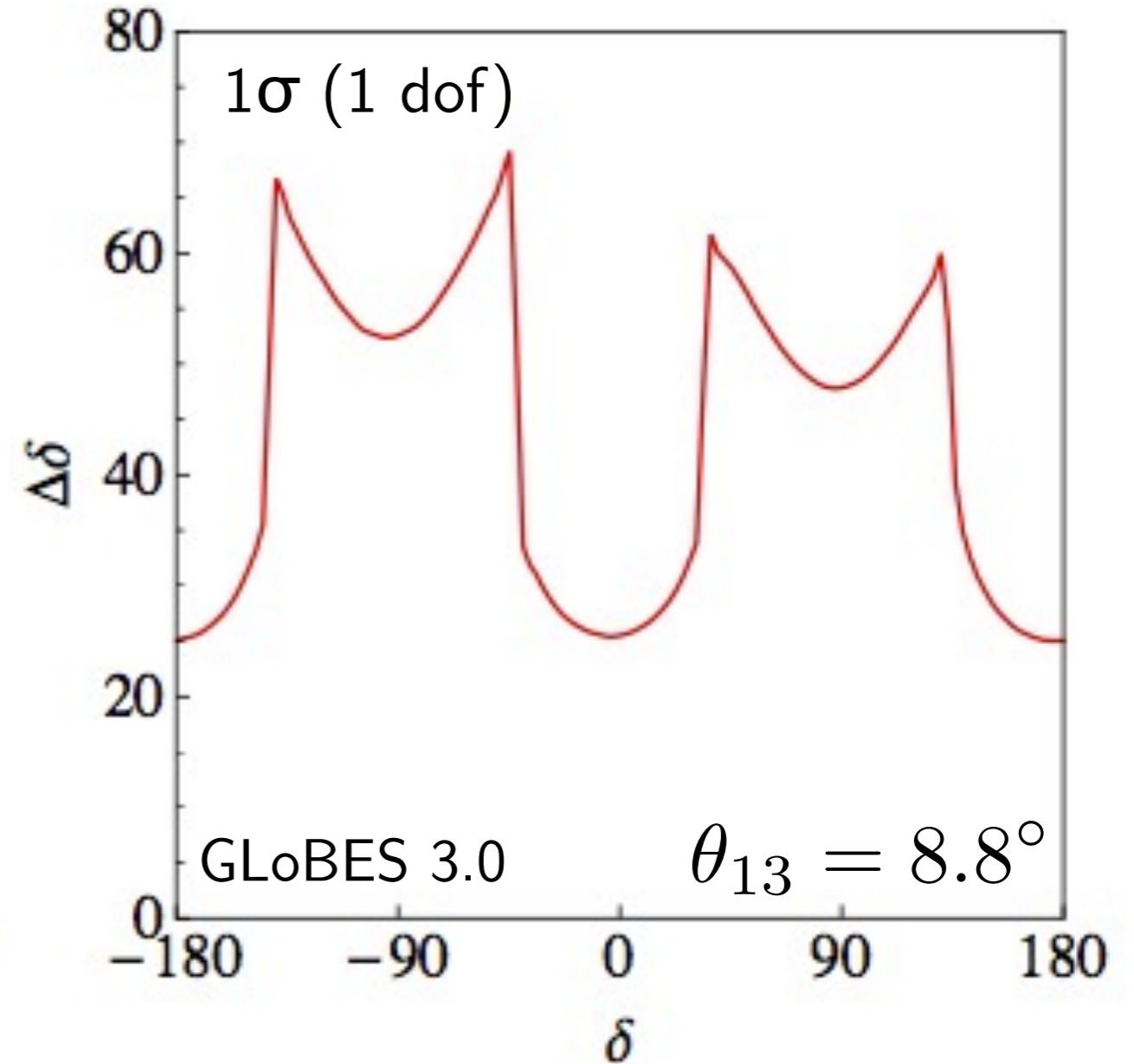
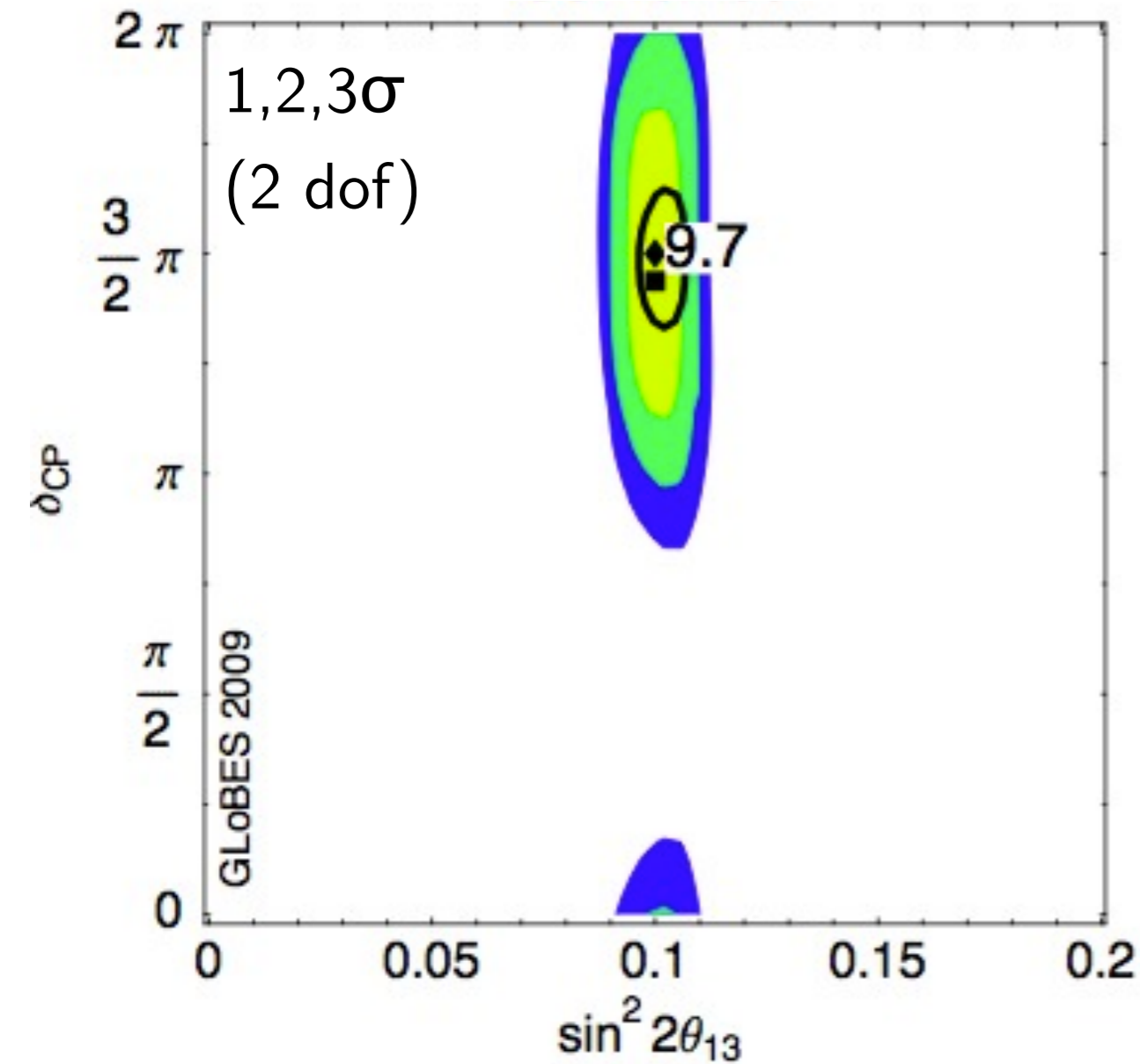
# Precision in $\delta$

- 1) Mild  $\theta_{13}$  dependence
- 2) Strong  $\delta$  dependence for BB350 due to no disappearance data
- 3) CPV discovery potential related to precision around  $0, \pi$ : more favorable for setups in vacuum and with similar number of  $\nu/\bar{\nu}$  events



# The starting point

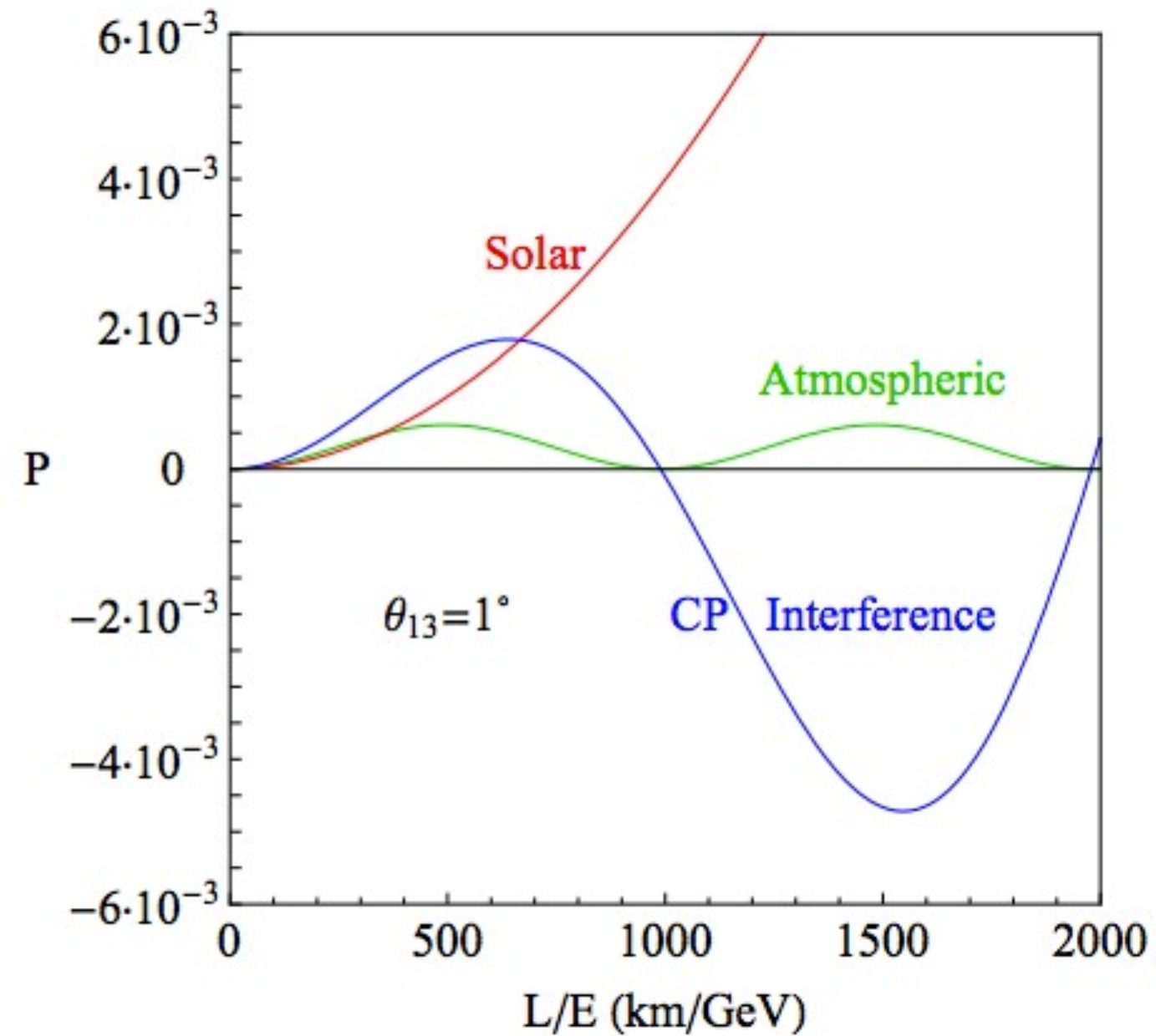
NO $\nu$ A+T2K+Daya Bay



Huber, Lindner, Schwetz, Winter,  
0907.1896 [hep-ph]

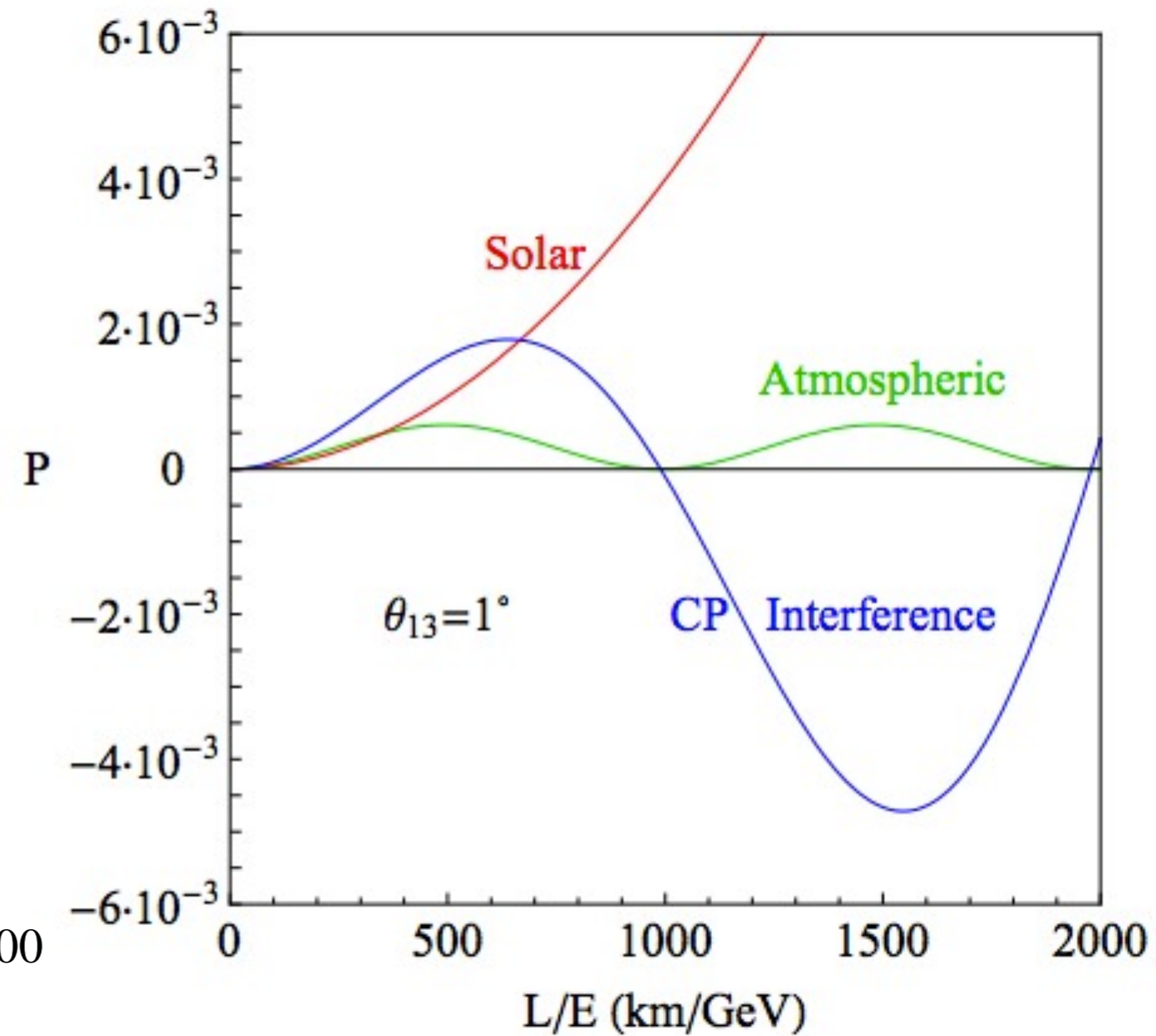
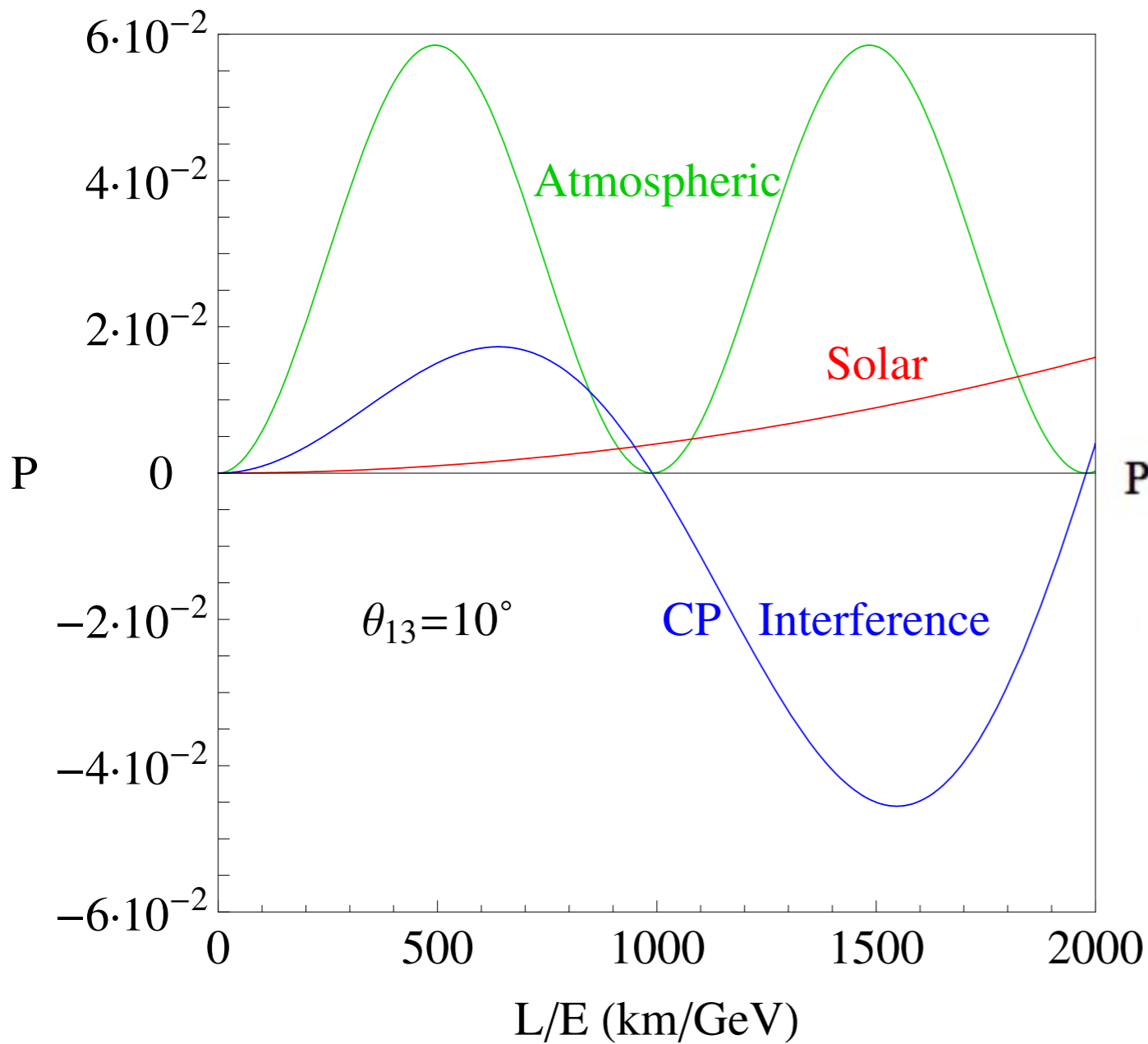
Coloma, Donini, Fernández-Martínez,  
Hernández, 1203.5651 [hep-ph]

# Impact of systematics



Coloma, Fernández-Martínez,  
1110.4583 [hep-ph]

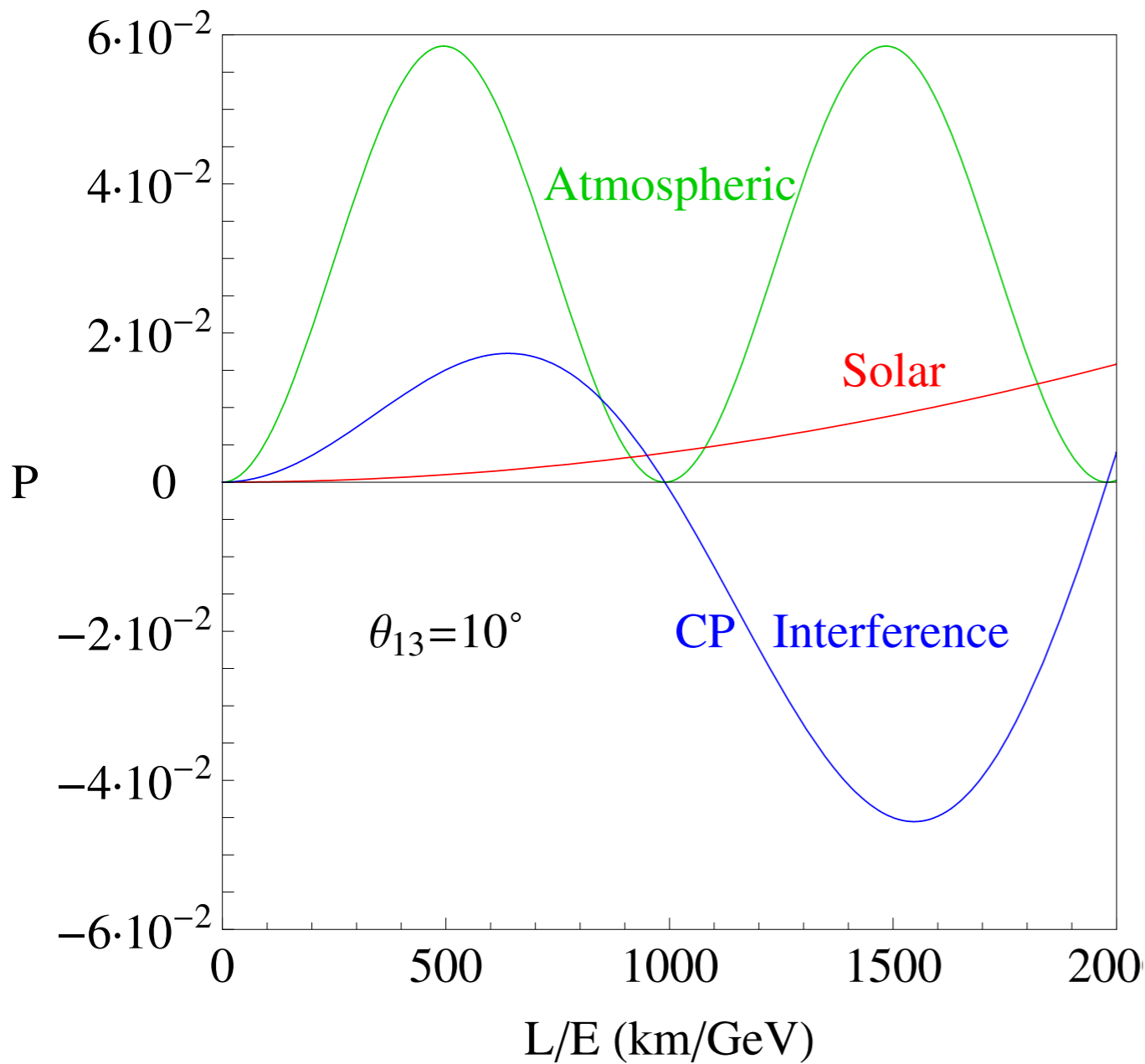
# Impact of systematics



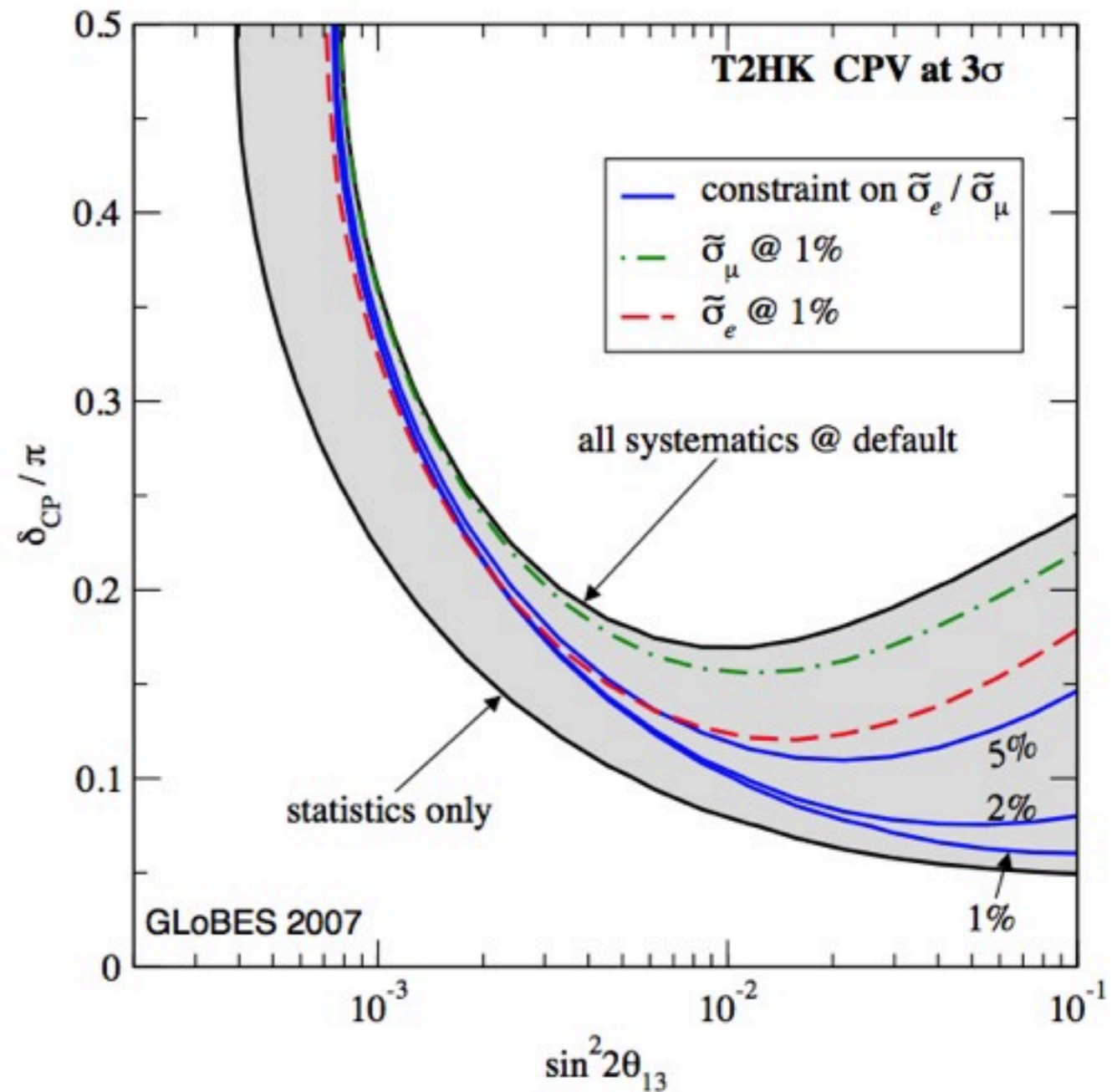
Coloma, Fernández-Martínez,  
1110.4583 [hep-ph]



# Impact of systematics



Coloma, Fernández-Martínez,  
1110.4583 [hep-ph]



Huber, Mezzetto, Schwetz,  
0711.2950 [hep-ph]

# An example

Possible ways to reduce the effect of systematics:

1) measure **final flavor cross sections** at a near detector.

If this cannot be done, put constraints on **ratios**

between cross sections for different flavors

Day, McFarland, 1206.6745 [hep-ph]

(see also Debbie's talk)

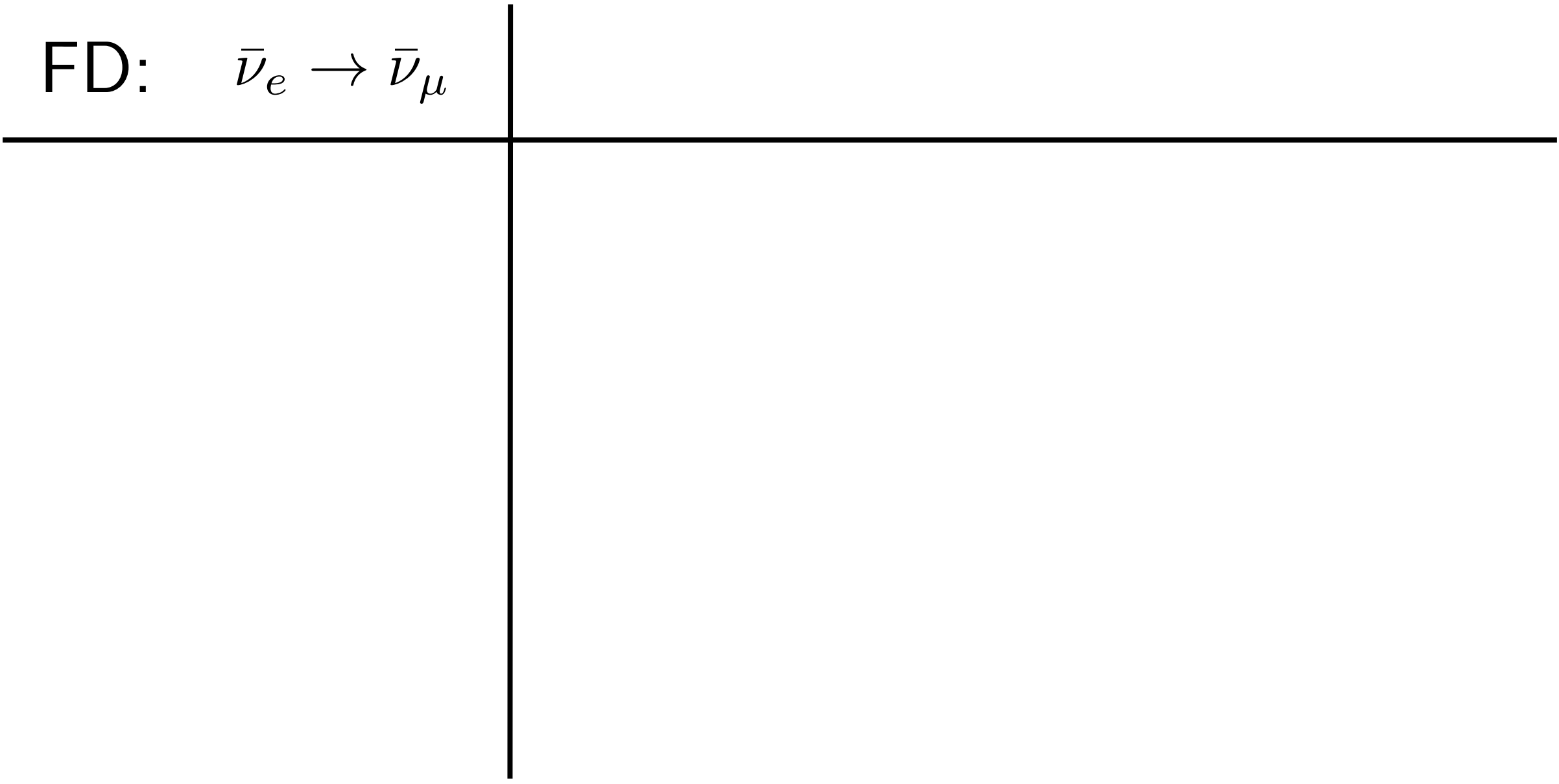
2) measure **intrinsic background** at near detector

3) use **data from disappearance** channels at the far detector

# An example

$$\mu^- \rightarrow e^- \bar{\nu}_e \nu_\mu$$

FD:  $\bar{\nu}_e \rightarrow \bar{\nu}_\mu$



# An example

$$\mu^- \rightarrow e^- \bar{\nu}_e \nu_\mu$$

FD:  $\bar{\nu}_e \rightarrow \bar{\nu}_\mu$

$\mu^-$

$\nu_{\text{far}}$

Matter

Xsec



# An example

$$\mu^- \rightarrow e^- \bar{\nu}_e \nu_\mu$$

FD:  $\bar{\nu}_e \rightarrow \bar{\nu}_\mu$

$\mu^-$

$\nu_{\text{far}}$

Matter

Xsec

ND:  $\nu_\mu \rightarrow \nu_\mu$

# An example

$$\mu^- \rightarrow e^- \bar{\nu}_e \nu_\mu$$

FD:	$\bar{\nu}_e \rightarrow \bar{\nu}_\mu$	Mu-	$\nu_{\text{far}}$	Matter	Xsec
ND:	$\nu_\mu \rightarrow \nu_\mu$	Mu-	$\nu_{\text{near}}$	Vac	Xsec

# An example

$$\mu^- \rightarrow e^- \bar{\nu}_e \nu_\mu$$

FD:	$\bar{\nu}_e \rightarrow \bar{\nu}_\mu$	Mu-	$V_{far}$	Matter	Xsec
ND:	$\nu_\mu \rightarrow \nu_\mu$	Mu-	$V_{near}$	Vac	Xsec
	$\bar{\nu}_\mu \rightarrow \bar{\nu}_\mu$	Mu+	$V_{near}$	Vac	Xsec

# An example

$$\mu^- \rightarrow e^- \bar{\nu}_e \nu_\mu$$

FD:	$\bar{\nu}_e \rightarrow \bar{\nu}_\mu$	Mu-	$V_{far}$	Matter	Xsec
ND:	$\nu_\mu \rightarrow \nu_\mu$	Mu-	$V_{near}$	Vac	Xsec
	$\bar{\nu}_\mu \rightarrow \bar{\nu}_\mu$	Mu+	$V_{near}$	Vac	Xsec
FD:	$\nu_\mu \rightarrow \nu_\mu$				
	$\bar{\nu}_\mu \rightarrow \bar{\nu}_\mu$				

# An example

$$\mu^- \rightarrow e^- \bar{\nu}_e \nu_\mu$$

FD:	$\bar{\nu}_e \rightarrow \bar{\nu}_\mu$	Mu-	$V_{far}$	Matter	Xsec
ND:	$\nu_\mu \rightarrow \nu_\mu$	Mu-	$V_{near}$	Vac	Xsec
	$\bar{\nu}_\mu \rightarrow \bar{\nu}_\mu$	Mu+	$V_{near}$	Vac	Xsec
FD:	$\nu_\mu \rightarrow \nu_\mu$	Mu-	$V_{far}$	Matter	Xsec
	$\bar{\nu}_\mu \rightarrow \bar{\nu}_\mu$	Mu+	$V_{far}$	Matter	Xsec

# Simulation details

Systematics	SB			BB			NF		
	Opt.	Def.	Cons.	Opt.	Def.	Cons.	Opt.	Def.	Cons.
Fiducial volume ND	0.2%	0.5%	1%	0.2%	0.5%	1%	0.2%	0.5%	1%
Fiducial volume FD (incl. near-far extrap.)	1%	2.5%	5%	1%	2.5%	5%	1%	2.5%	5%
Flux error signal $\nu$	5%	7.5%	10%	1%	2%	2.5%	0.1%	0.5%	1%
Flux error background $\nu$	10%	15%	20%	correlated			correlated		
Flux error signal $\bar{\nu}$	10%	15%	20%	1%	2%	2.5%	0.1%	0.5%	1%
Flux error background $\bar{\nu}$	20%	30%	40%	correlated			correlated		
Background uncertainty	5%	7.5%	10%	5%	7.5%	10%	10%	15%	20%
Cross secs $\times$ eff. QE <sup>†</sup>	10%	15%	20%	10%	15%	20%	10%	15%	20%
Cross secs $\times$ eff. RES <sup>†</sup>	10%	15%	20%	10%	15%	20%	10%	15%	20%
Cross secs $\times$ eff. DIS <sup>†</sup>	5%	7.5%	10%	5%	7.5%	10%	5%	7.5%	10%
Effec. ratio $\nu_e/\nu_\mu$ QE <sup>*</sup>	3.5%	11%	–	3.5%	11%	–	–	–	–
Effec. ratio $\nu_e/\nu_\mu$ RES <sup>*</sup>	2.7%	5.4%	–	2.7%	5.4%	–	–	–	–
Effec. ratio $\nu_e/\nu_\mu$ DIS <sup>*</sup>	2.5%	5.1%	–	2.5%	5.1%	–	–	–	–
Matter density	1%	2%	5%	1%	2%	5%	1%	2%	5%

(details in 1209.5973 [hep-ph])



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(details in 1209.5973 [hep-ph])



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Effec. ratio $\nu_e/\nu_\mu$ QE*	3.5%	11%	–	3.5%	11%	–	–	–	–
Effec. ratio $\nu_e/\nu_\mu$ RES*	2.7%	5.4%	–	2.7%	5.4%	–	–	–	–
Effec. ratio $\nu_e/\nu_\mu$ DIS*	2.5%	5.1%	–	2.5%	5.1%	–	–	–	–
Matter density	1%	2%	5%	1%	2%	5%	1%	2%	5%

theoretical constraint

(details in 1209.5973 [hep-ph])



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Matter density	1%	2%	5%	1%	2%	5%	1%	2%	5%

theoretical constraint

(details in 1209.5973 [hep-ph])

# Simulation details

$$\chi^2 = \sum_{r,i} 2 \left( T_{r,i}(\vec{\Theta}, \vec{\xi}) - O_{r,i} + O_{r,i} \ln \frac{O_{r,i}}{T_{r,i}(\vec{\Theta}, \vec{\xi})} \right) + \sum_k \left( \frac{\xi_k}{\sigma_k} \right)^2$$

- GLoBES software used hep-ph/0407333, hep-ph/0701187
- A near detector has been explicitly simulated for all experiments
- Correlations are fully taken into account between different channels (unless otherwise stated)
- Systematic uncertainties introduced as nuisance parameters
- Marginalization performed over all parameters
- No degeneracies considered. Normal hierarchy assumed.
- $\sin^2 2\theta_{13} = 0.1$

# More simulation details

- No energy dependent effects included (nuclear effects, for instance)

see Zeller's and Morfin's talks

All errors included as norm errors.

- However, different independent errors considered for the different cross section regimes: “effective” shape error
- Near detector is assumed to be sufficiently far away so that the spectrum is identical to the far detector (1-2 km)
- Near and far detector are assumed to be identical (except for the treatment of NC backgrounds)
- No tau backgrounds are included (very little impact for CPV, though)

Donini, Gomez Cadenas, Meloni, 1005.2275 [hep-ph]

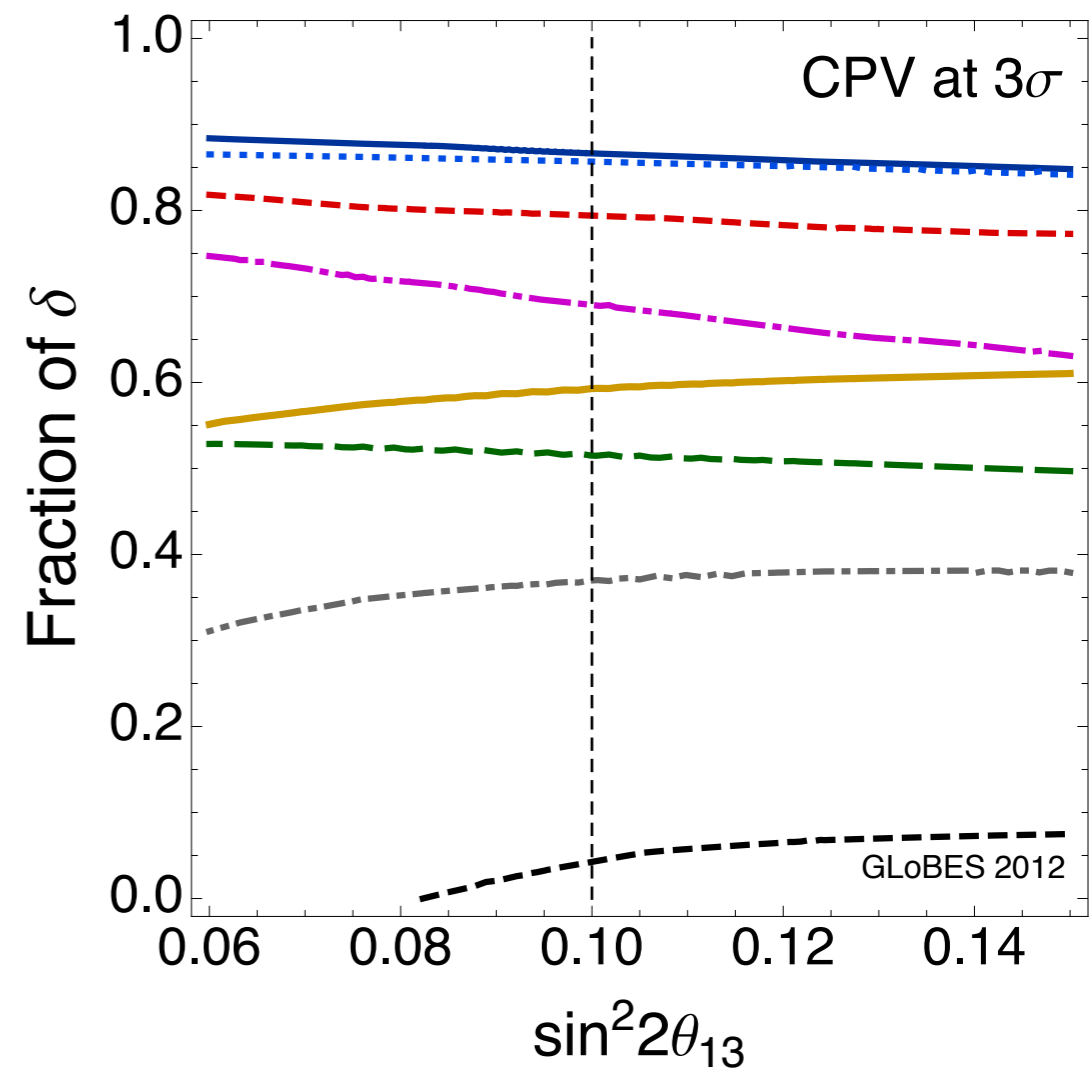
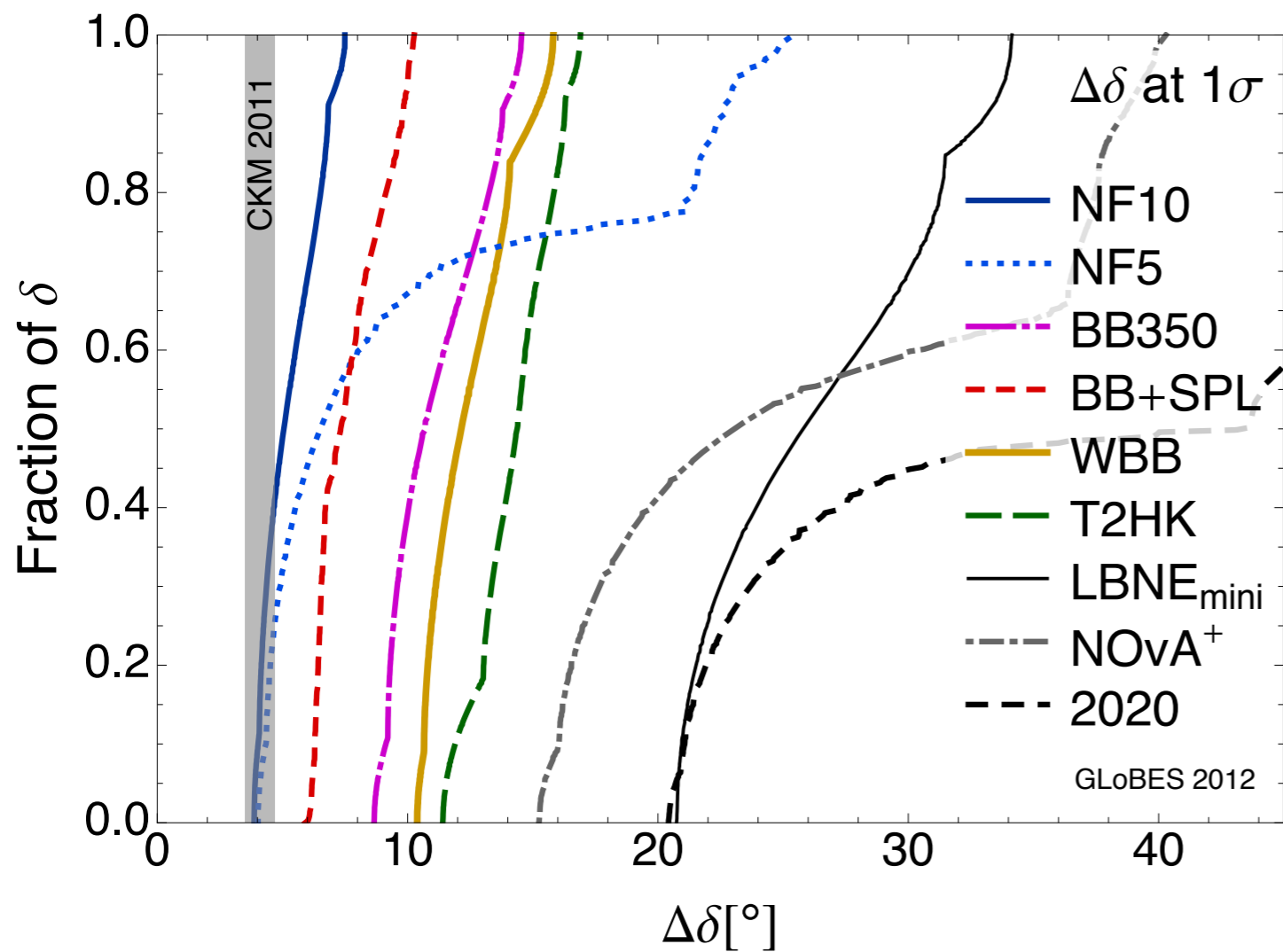


# The setups

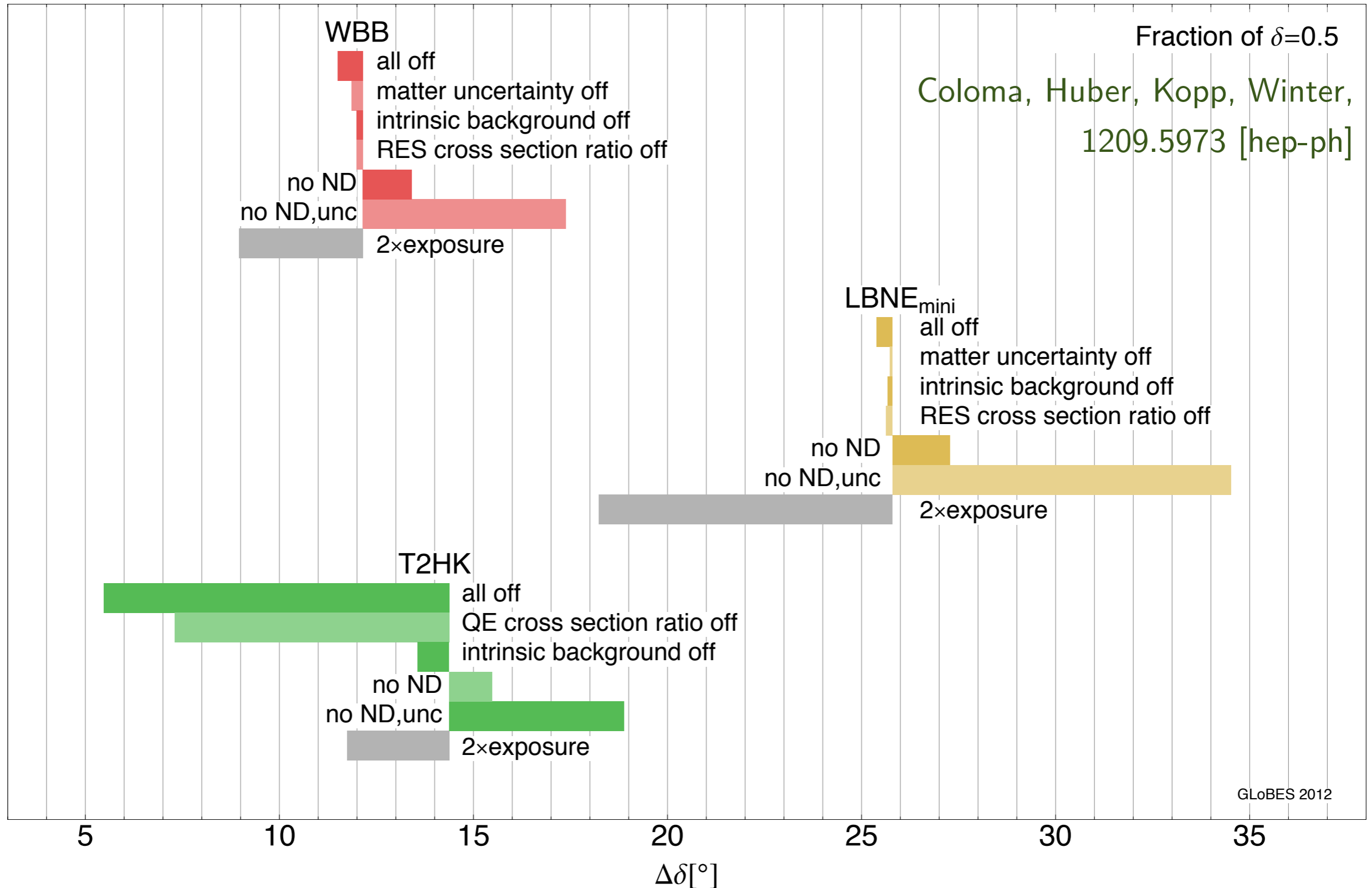
	Setup	$E_\nu^{\text{peak}}$	$L$	OA	Detector	kt	MW	Decays/yr	$(t_\nu, t_{\bar{\nu}})$
Benchmark	BB350	1.2	650	–	WC	500	–	$1.1(2.8) \times 10^{18}$	(5,5)
	NF10	5.0	2000	–	MIND	100	–	$7 \times 10^{20}$	(10,10)
	WBB	4.5	2300	–	LAr	100	0.8	–	(5,5)
	T2HK	0.6	295	$2.5^\circ$	WC	560	1.66	–	(1.5,3.5)
Alternative	BB100	0.3	130	–	WC	500	–	$1.1(2.8) \times 10^{18}$	(5,5)
	+ SPL			–			4		–
	NF5	2.5	1290	–	MIND	100	–	$7 \times 10^{20}$	(10,10)
	LBNE <sub>mini</sub>	4.0	1290	–	LAr	10	0.7	–	(5,5)
	NO $\nu$ A <sup>+</sup>	2.0	810	$0.8^\circ$	LAr	30	0.7	–	(5,5)
2020	T2K	0.6	295	$2.5^\circ$	WC	22.5	0.75	–	(5,5)
	NO $\nu$ A	2.0	810	$0.8^\circ$	TASD	15	0.7	–	(4,4)

# General comparison

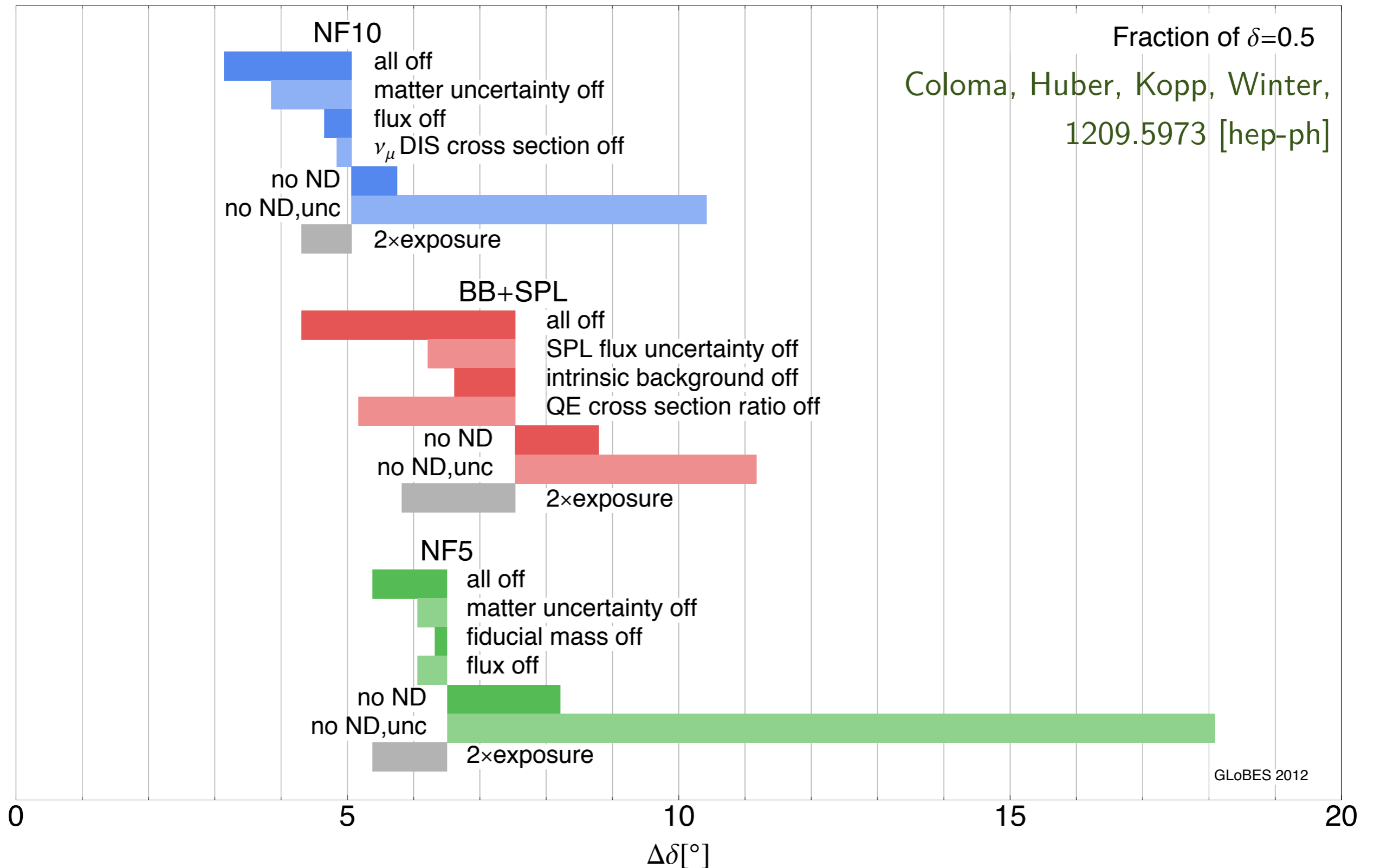
How **far** do we want to get?



# Sys and near detectors

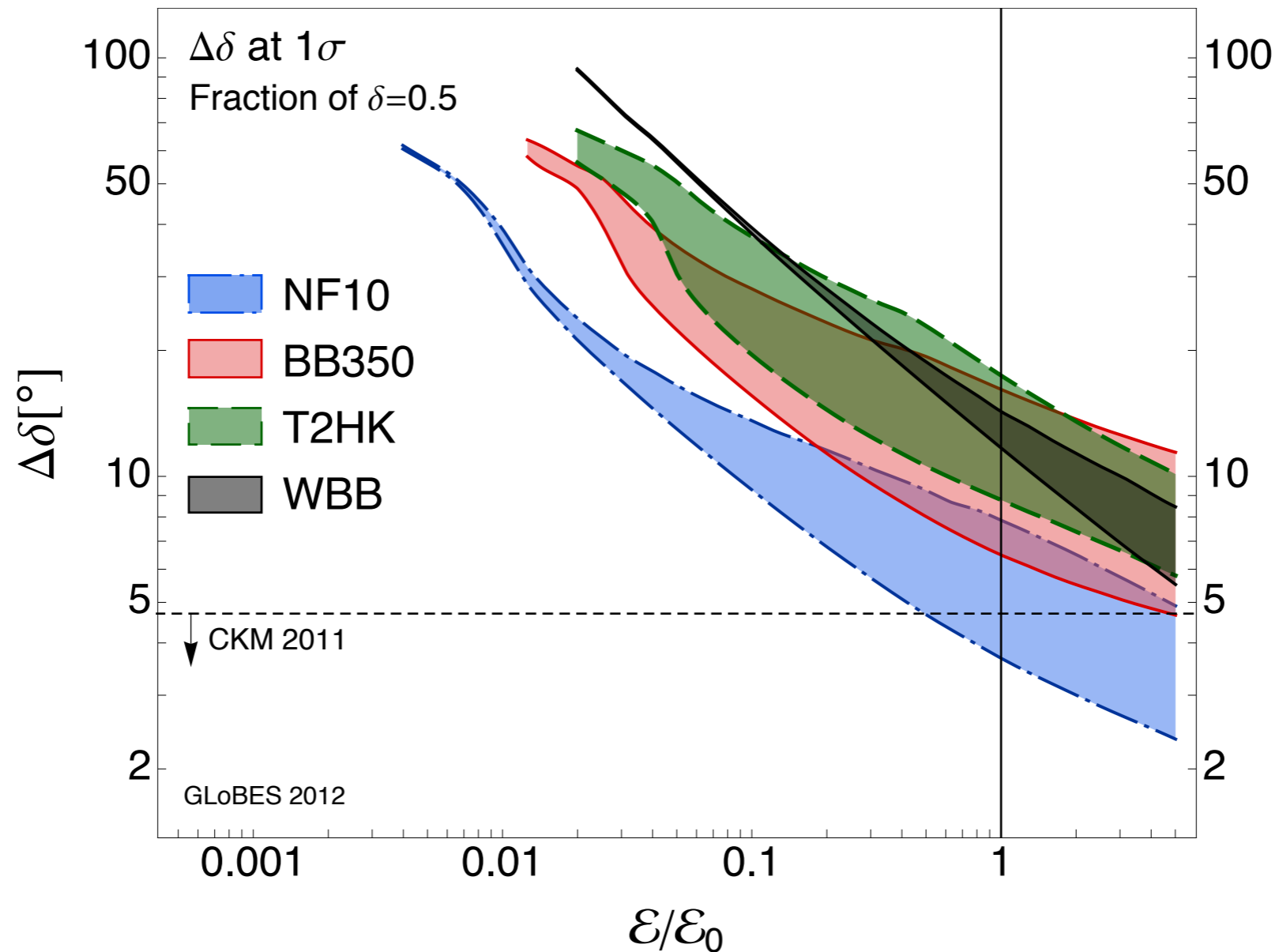


# Sys and near detectors



# Exposure vs systematics

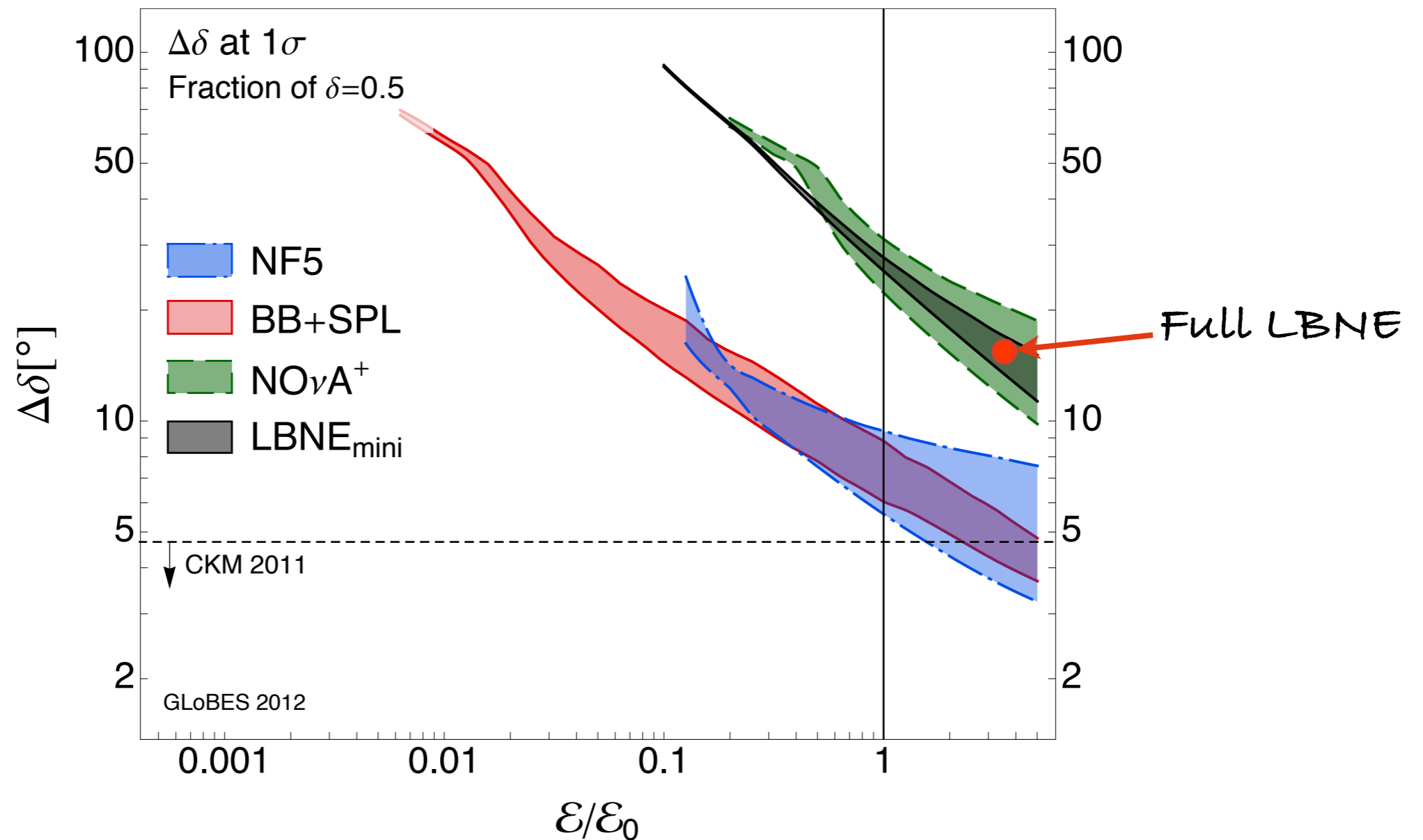
Variation between optimistic and conservative assumptions:





# Exposure vs systematics

Variation between optimistic and conservative assumptions:



# Conclusions

- The precision on  $\theta_{13}$  obtained at Daya Bay will most likely not be exceeded by any beam experiment. For  $\delta$  the situation is more complicated, though:
  - ideally, an experiment in vacuum would have the best reach for CPV; however, this is not optimal for precision...
  - maybe a combination of the two?
- Low energy setups are generally more affected by systematics
  - theoretical assumptions on cross section ratios are critical (Exception! BB+SPL)

# Conclusions

- The impact of a ND does not seem so relevant if data from disappearance at the FD is used (under certain assumptions!)
  - migration due to nuclear effects is not included
  - if NP is present, a ND is crucial
  - we assume only norm errors (an “effective” shape error is done for the xsecs only)
  - effectiveness of this method depends on the statistics
- In some cases, it may a be better path to increase statistics than reduce systematics...
- LENF is the only facility able to achieve similar precision to quark sector

Backup

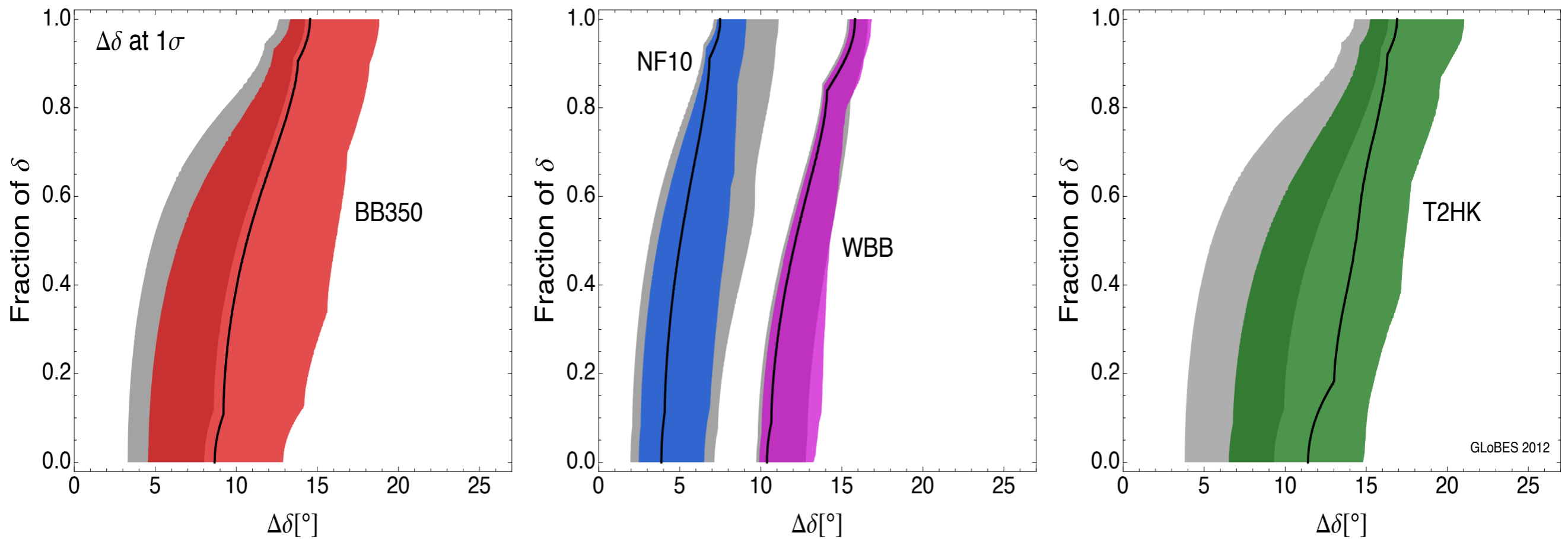
# Impact of systematics

Systematics become a problem for large  $q_{13}$  since the leading term in the probability grows quadratically

$$P_{e\mu}^{\pm} = X_{\pm} \sin^2 2\theta_{13} + Z + Y_{\pm} \cos \theta_{13} \sin 2\theta_{13} \cos \left( \pm\delta - \frac{\Delta_{31}L}{2} \right)$$

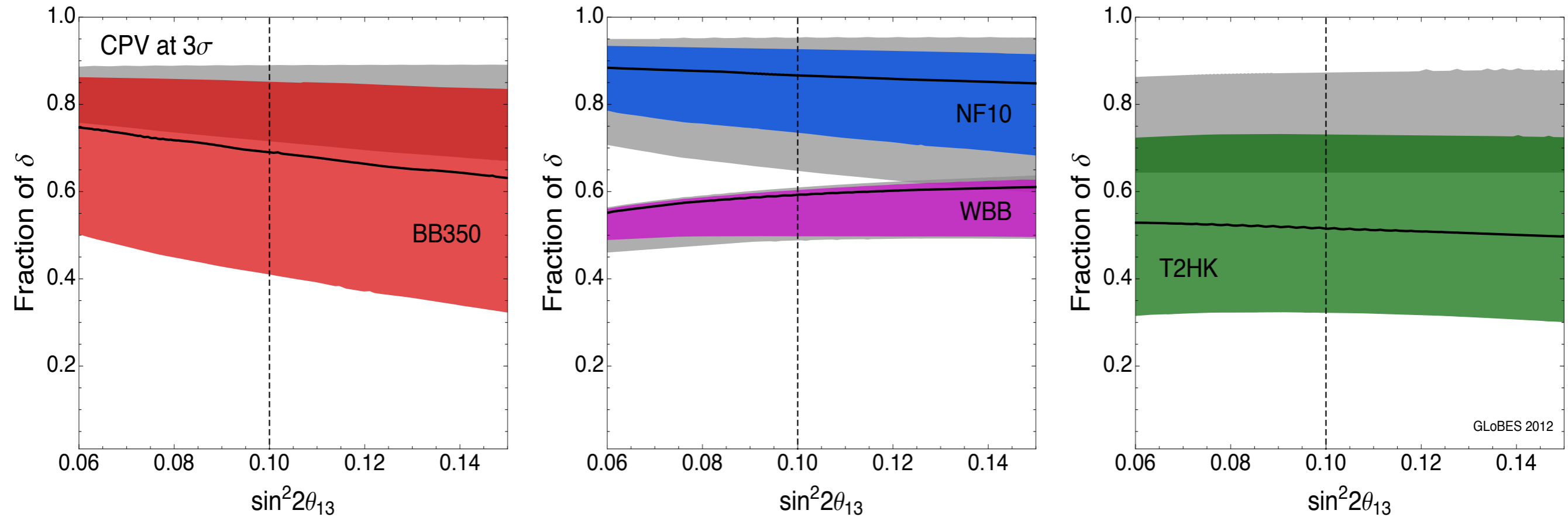
# Impact of systematics

Variation between optimistic and conservative assumptions:



# Impact of systematics

Variation between optimistic and conservative assumptions:



# Mass hierarchy

Mass hierarchy may be obtained through:

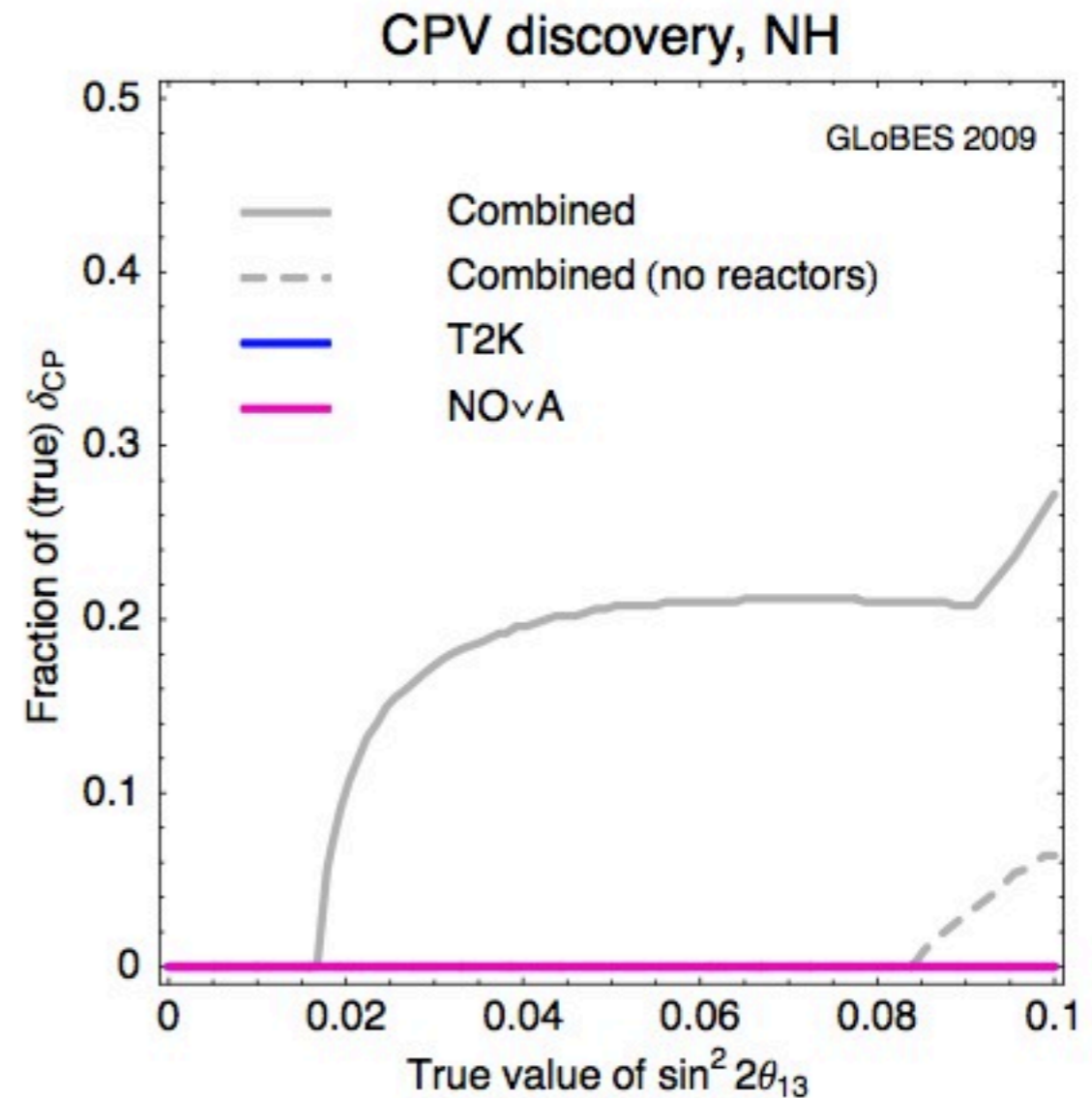
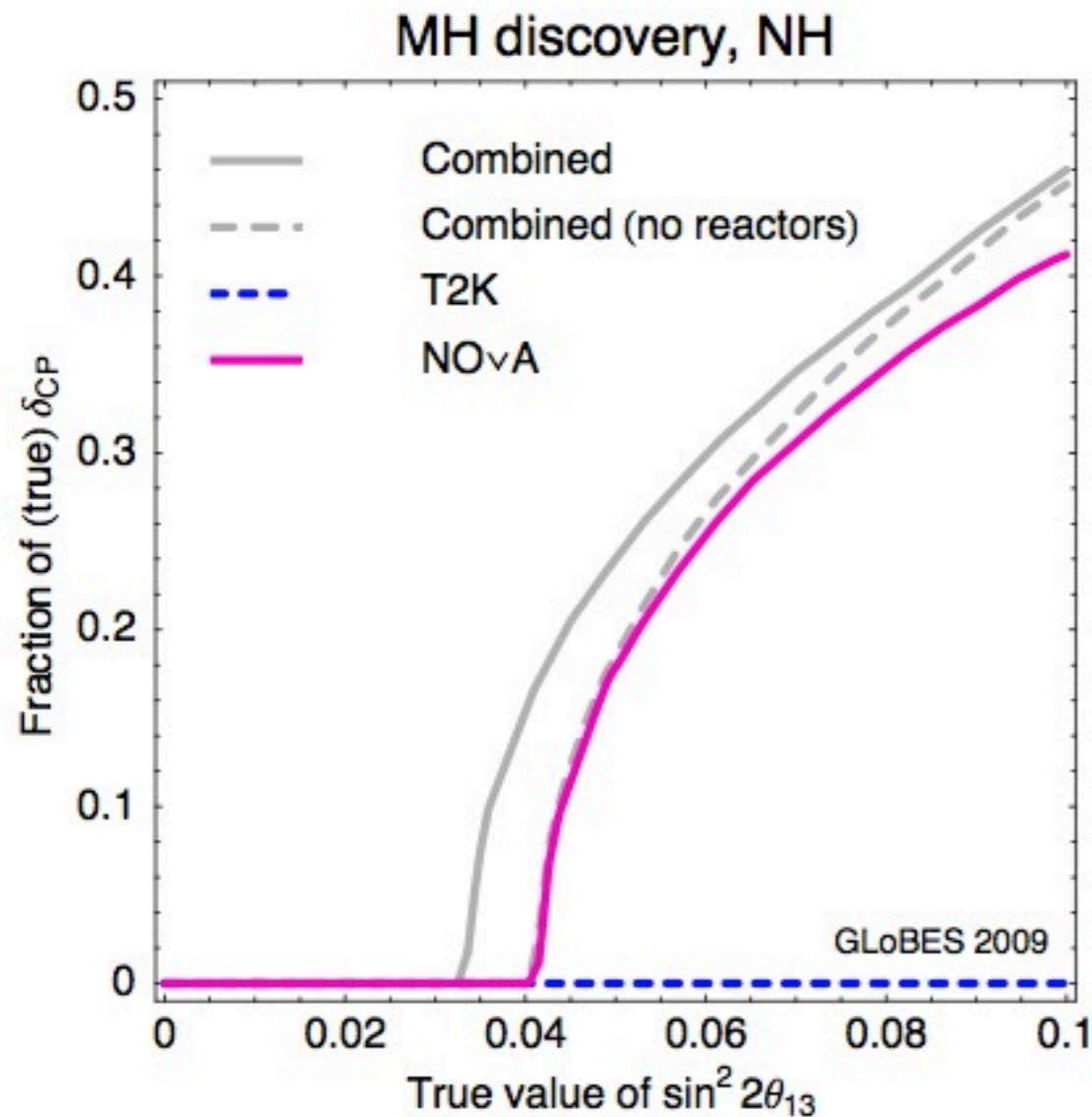
- T2K+NOvA+INO 1203.3388 [hep-ph]
- Atmospheric data at future exps 1109.3262 [hep-ex]
- PINGU 1205.7071 [hep-ph]
- Daya Bay II hep-ph/0112074
- combination of precise reactor+LBL data  
hep-ph/0503283
- ...

...or a combination of all of them!



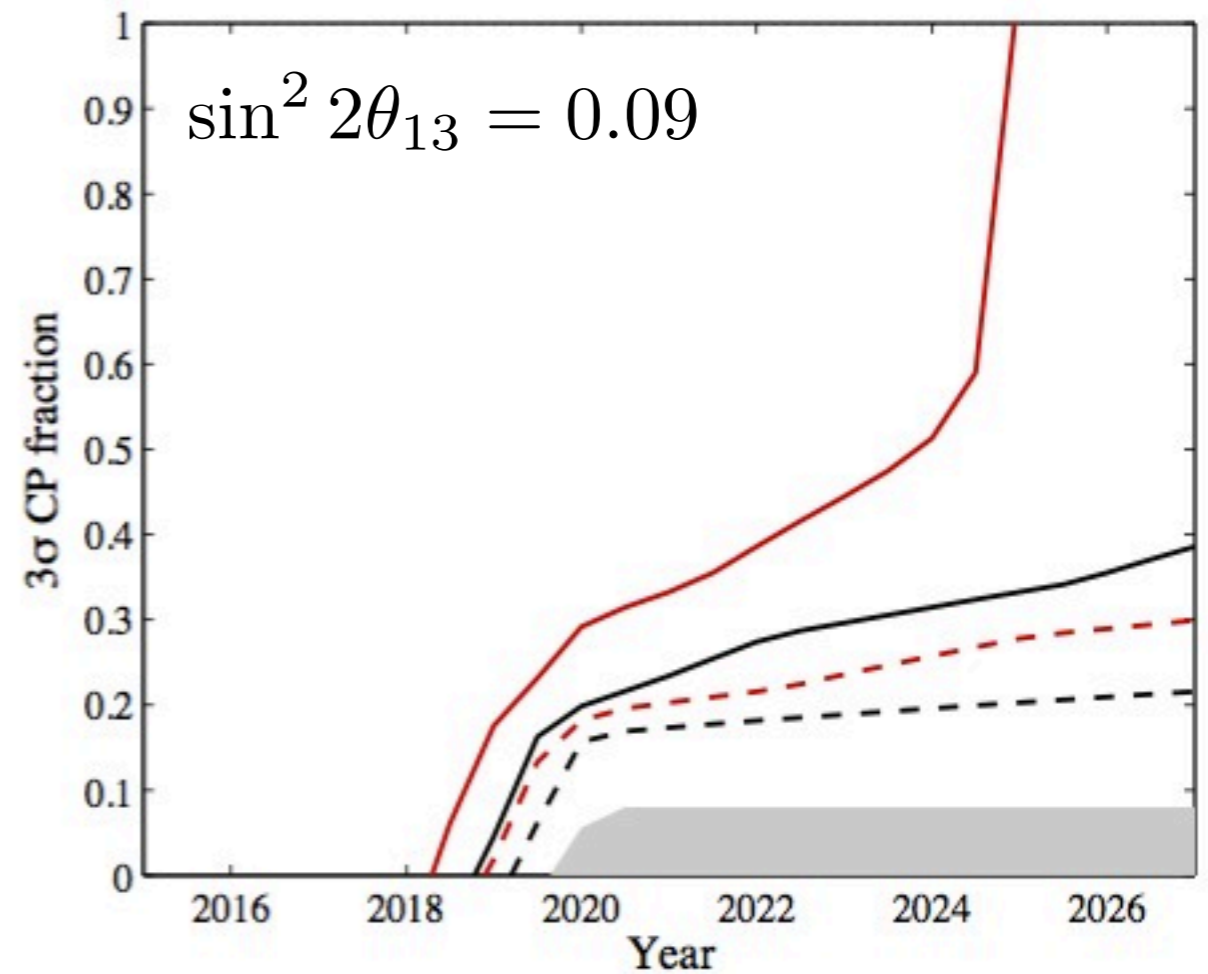
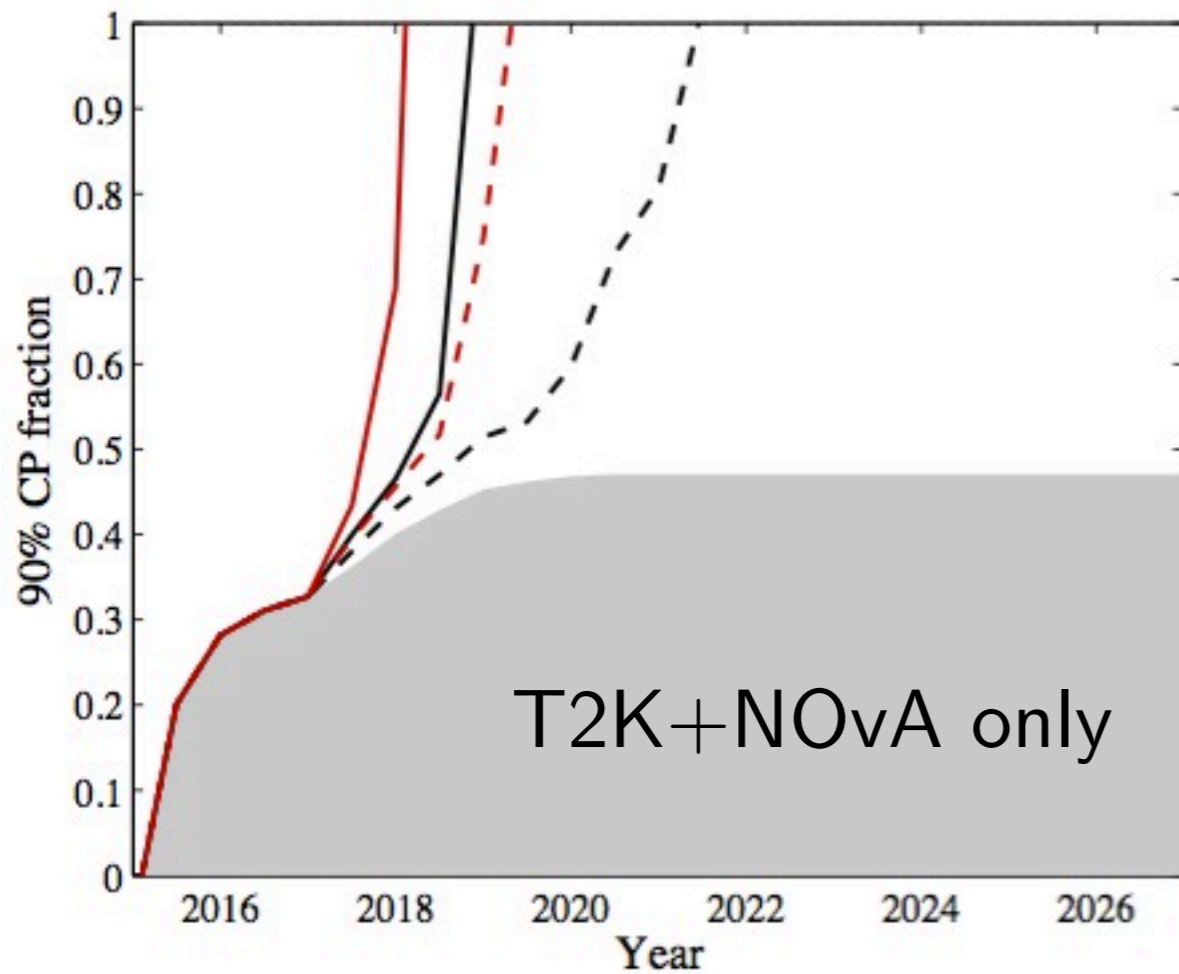
# Present oscillation facilities

Discovery potential at the 90% CL



# Present oscillation facilities

T2K+NOvA+INO  
(50kt/100kt; low/high res)



Blennow, Schwetz, 1203.3388 [hep-ph]

# Previous hints on $\theta_{13}$

Previous hints from global fits pointed to nonzero  $\theta_{13}$ ...

